

Utah State University

From the Selected Works of David Smith

Spring 2017

High-frequency radio wave absorption in the D-region

David Alan Smith, *Utah State University*



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High-frequency Radio Wave Absorption

A research report presented by
David Alan Smith

Utah State University
Department of Physics
March 30, 2017

Game Plan

Radio Wave Propagation

- Sky Waves
- Properties of Ionosphere
- Geometric Optics

High-frequency Radio Wave Absorption

- Basic Absorption Equation
- Types of Absorption
- Absorption Coefficient
- Absorption Equation
- Special Cases

Electron Density

Collision Frequency

Conclusions/Discussion

Questions

Essential Points

Absorption is **frequency-dependent**

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Most HF absorption takes place within the **D-region**

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Within the D-region **non-deviative** absorption dominates
(Assuming HF)

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Absorption is frequency-dependent

Most HF absorption takes place within the D-region

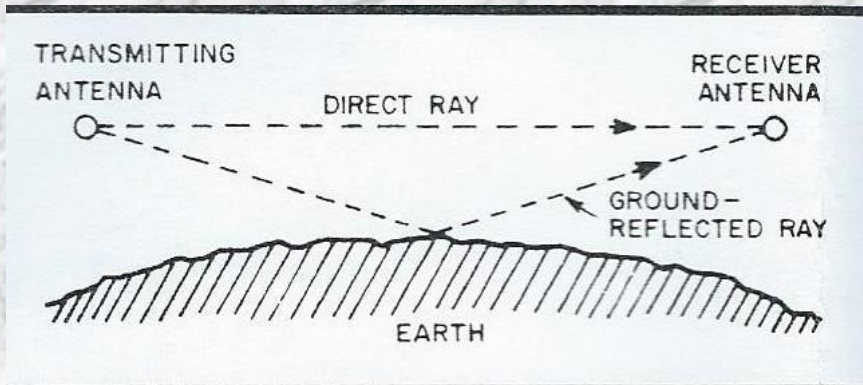
Within the D-region non-deviative absorption dominates
(Assuming HF)

Electron density is critical

Propagation via Sky Wave

Ground Wave vs. Sky Wave

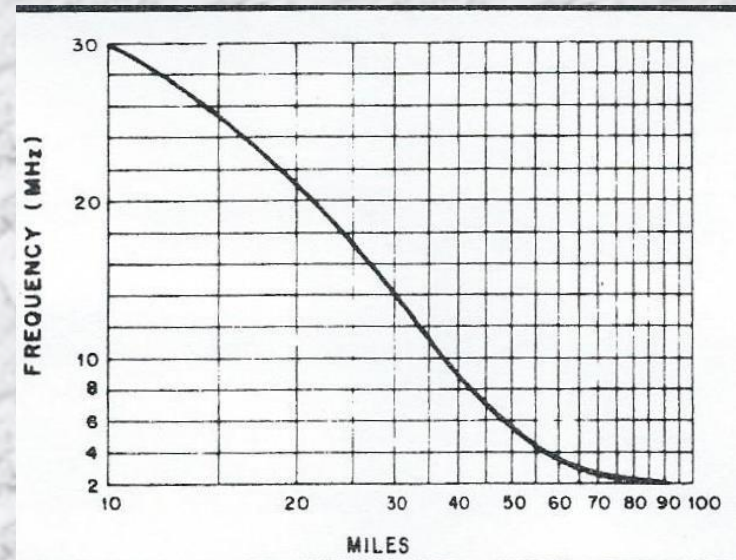
Ground Wave



Direct ray and ground-reflected ray
combine to form space wave

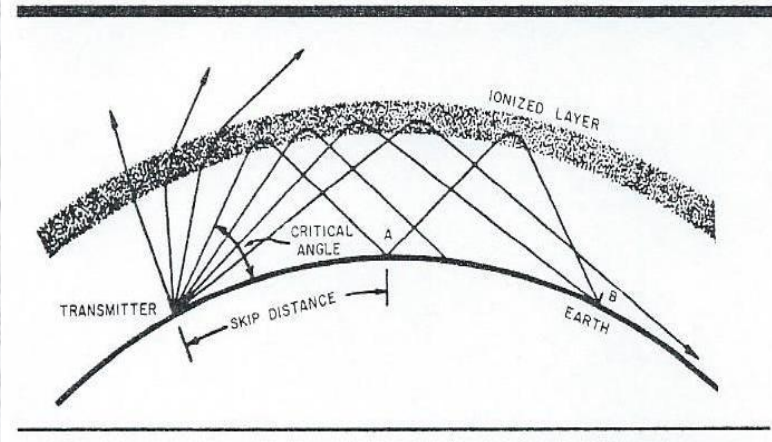
- AM broadcast band has range of about 160 km (100 miles)
- **Generally ineffective for long-range communications**

- Stays close to the earth
- Doesn't leave lower atmosphere



Typical HF ground-wave range

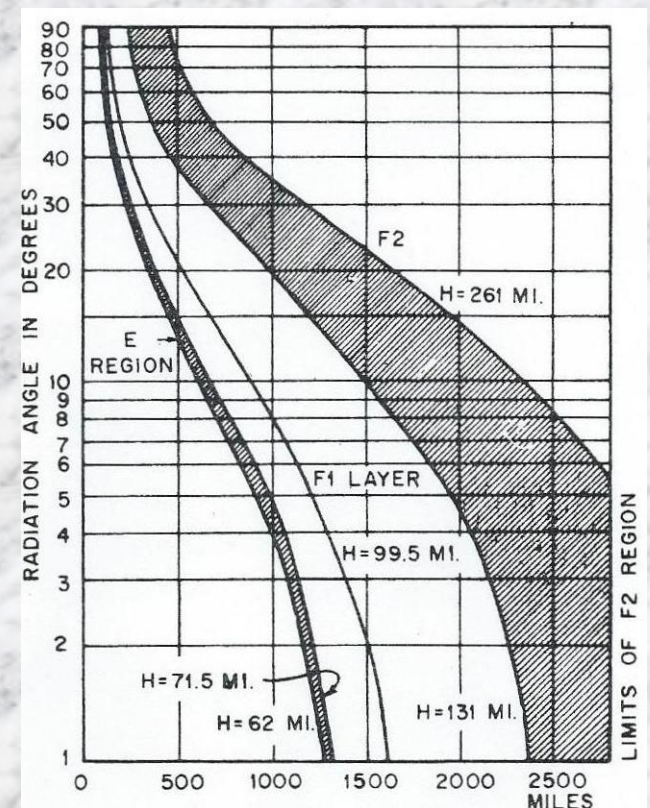
Sky Wave



Radio waves entering ionosphere at angles above critical angle go off into space

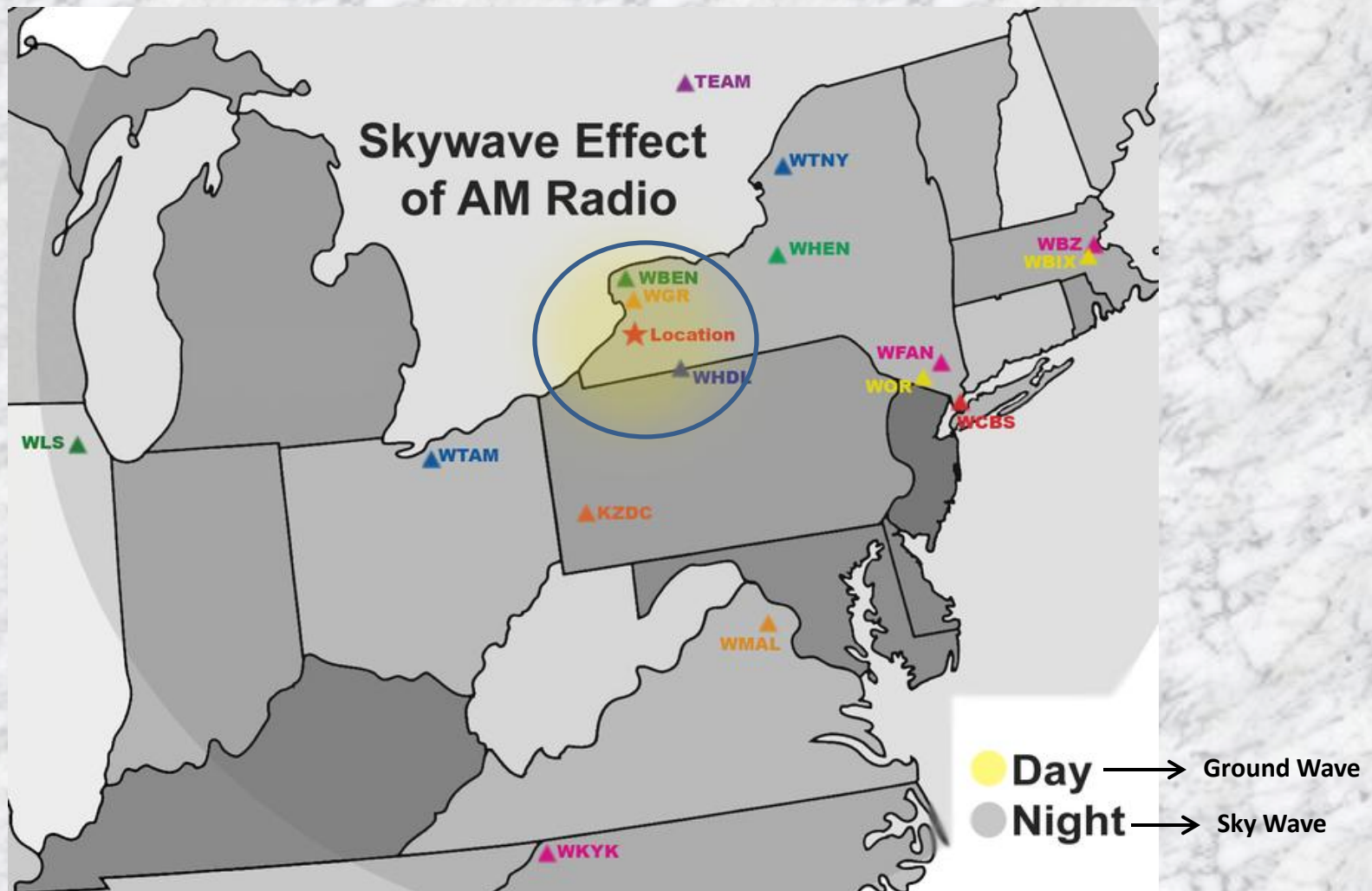
- Range up to 4000 km (2500 miles) per hop
- **Efficient long-range communication**
- *Subject to various atmospheric conditions*

- Leaves lower atmosphere
- Passes through ionized region
- Refracted according to geometric optics



Distance vs. wave angle for one-hop transmission

Ground Wave vs. Sky Wave



Definitions

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High-frequency Spectrum: 3.0 MHz-30 MHz

Ionospheric Properties

Ionosphere

The ionosphere is considered a weakly-ionized plasma

For a **fully-ionized** plasma the ratio of charged particles to neutral particles is about 1

Within the ionized region of the atmosphere this ratio is always much less than 1.

Hence the ionosphere is a weakly-ionized plasma

D-region

Height: About 90 km

Thickness: About 40 km

Significant diurnal variations

Typical daytime electron density

$$\sim 10^9 - 10^{11} \text{ m}^{-3}$$

E-region

Height: About 150 km

Thickness: About 60 km

Diurnal variations though not as pronounced as D-region

Typical daytime electron density:

$$\sim 1 \times 10^{11} - 4 \times 10^{11} \text{ m}^{-3}$$

F1-region

Height: About 350 km

Thickness: About 200 km

Diurnal variations

Typical daytime electron density:

$$\sim 4 \times 10^{11} - 2 \times 10^{12} \text{ m}^{-3}$$

F2-region

Height: About 1000 km

Thickness: About 750 km

Diurnal variations, though not as pronounced

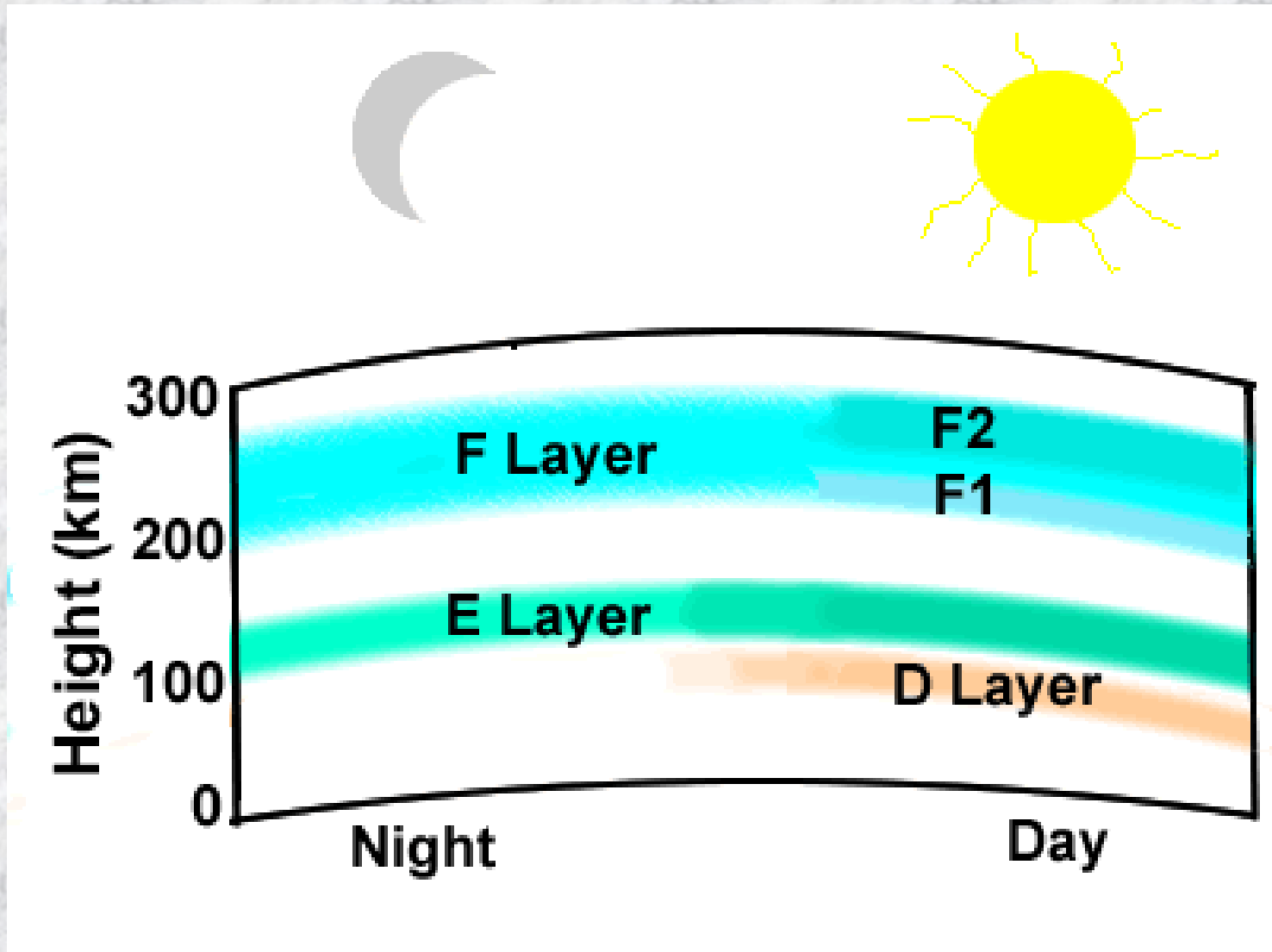
Electron Density:

$$\sim 8 \times 10^{10} - 2 \times 10^{12} \text{ m}^{-3}$$

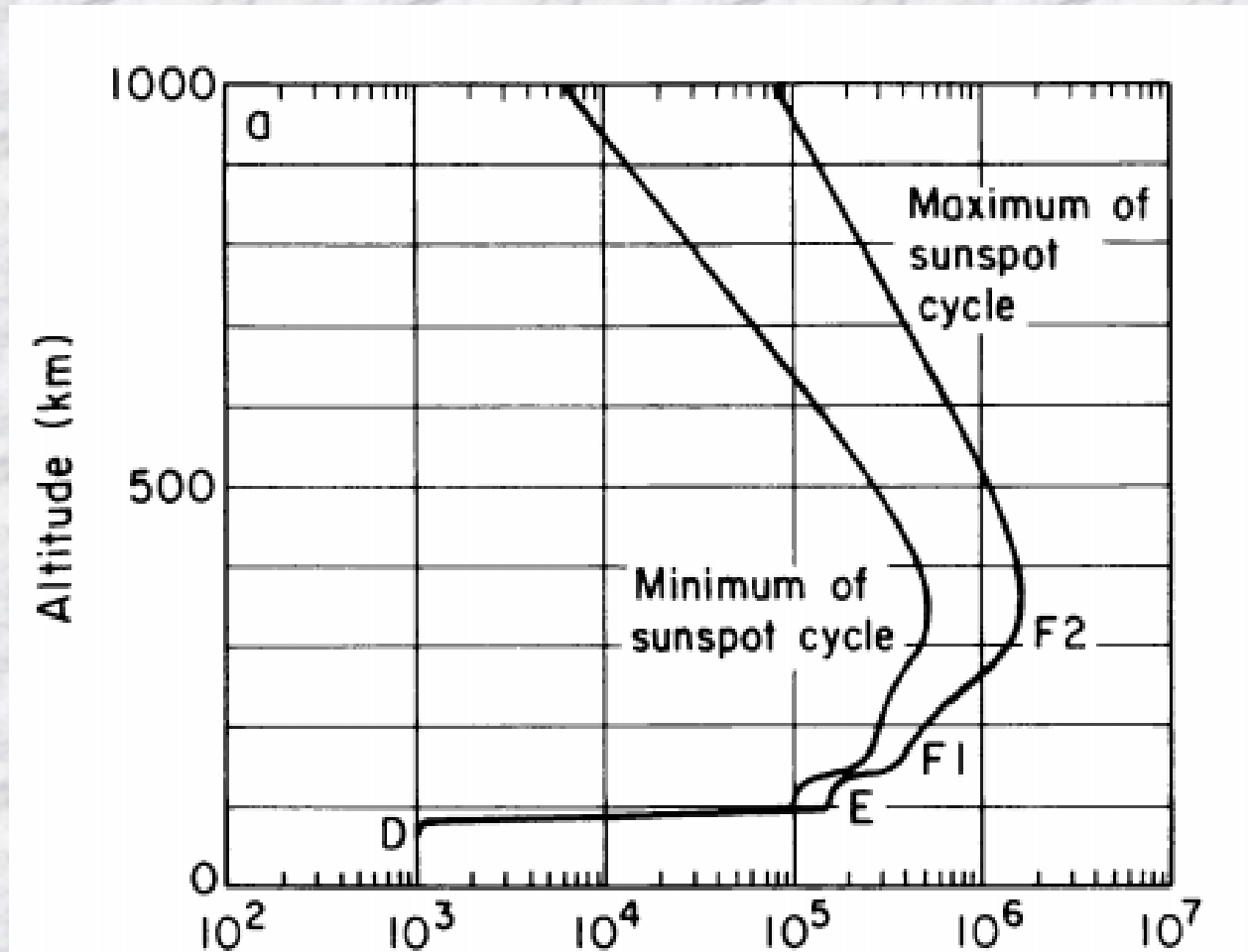
Unlike previous regions, F2 electron density decreases with height

Important note: F2 becomes the F-region after sunset.

Ionospheric Properties



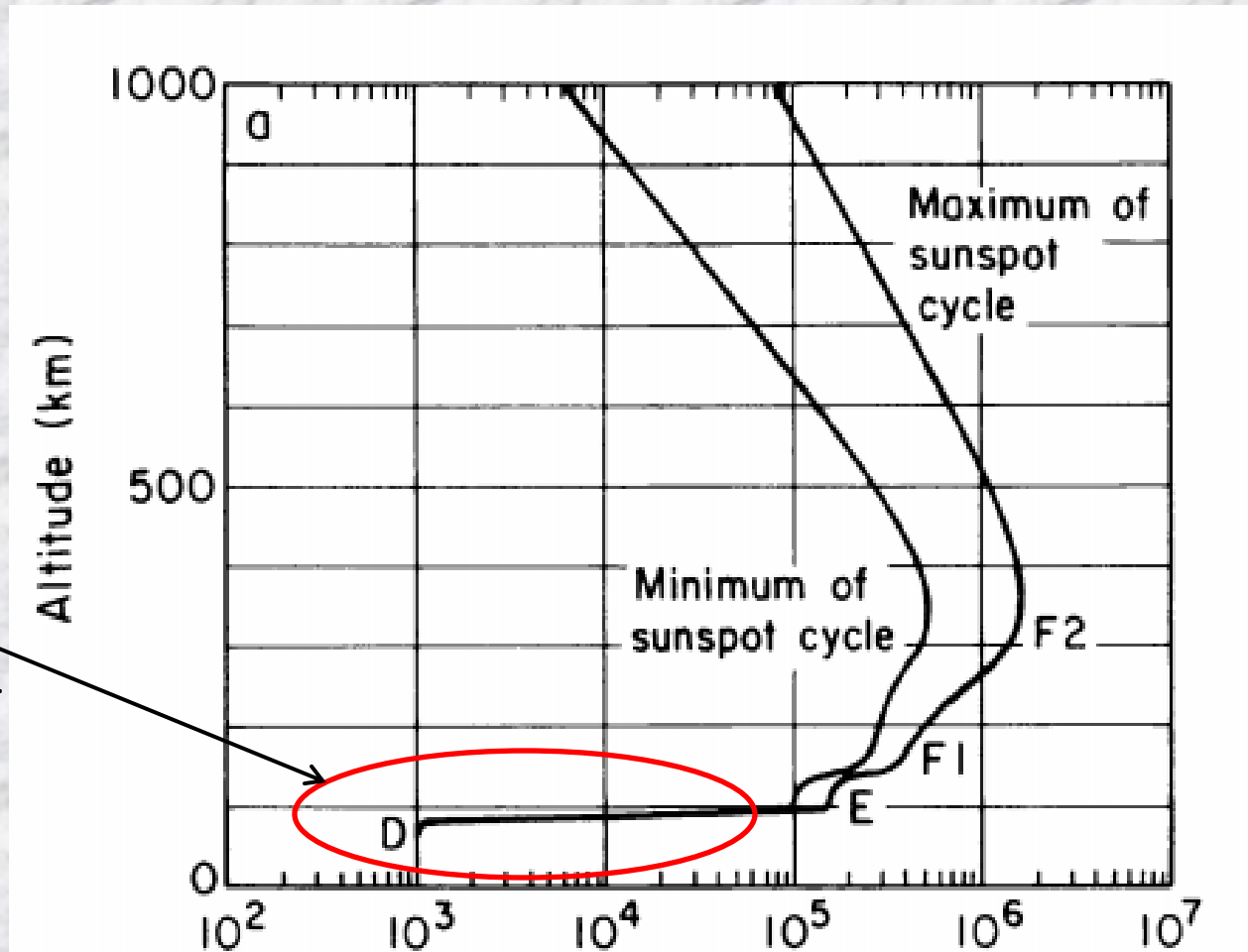
Electron Density: Function of Height



Electron concentration per cubic centimeter (Daytime)

Image from Kelley p 460

Electron Density: Function of Height

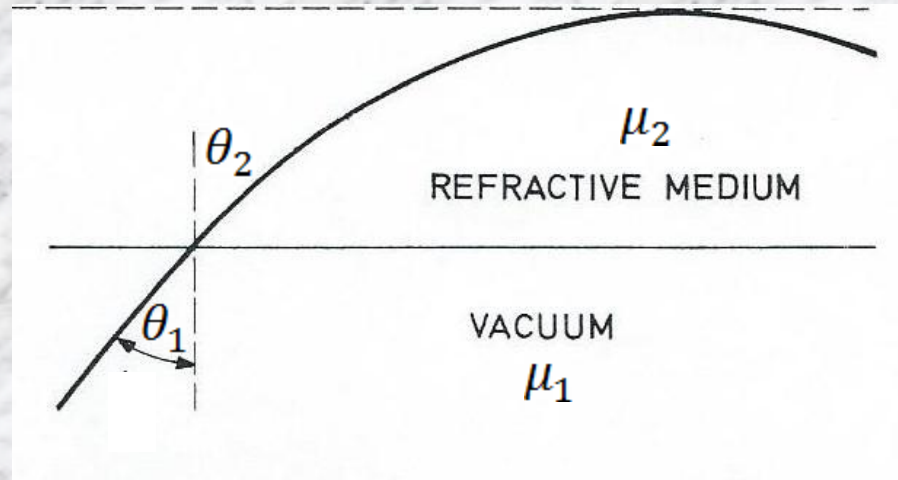


Electron concentration per cubic centimeter (Daytime)

Image from Kelley p 460

Geometric Approach

Geometric Approach



Ray path in a continuously varying medium (Ionosphere)
(Lied p 4)

When $\mu_2 < \mu_1 \Rightarrow \theta_1 < \theta_2$



“Bends away from the normal”

Geometric Approach

Snell's Law:

$$\mu_1 \sin(\theta_1) = \mu_2 \sin(\theta_2)$$

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Index of Refraction: $\mu = \frac{c}{v_{phase}}$

$$\mu = \sqrt{\epsilon_r}$$

Geometric Approach

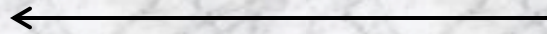
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Dielectric Constant of Weakly-Ionized Gas

$$\epsilon_r = 1 - \frac{\omega_p^2}{\omega^2}$$



$$\mu = \sqrt{\epsilon_r}$$



Geometric Approach

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$$\omega = 2\pi f$$

$$\mu = \sqrt{1 - \frac{n_e e^2}{4\pi^2 \epsilon_0 m_e f^2}}$$

Note dependence on **electron density**

Geometric Approach

$$\mu = \sqrt{1 - \frac{n_e e^2}{4\pi^2 \epsilon_0 m f^2}}$$

Within a dispersive media such as the ionosphere:

Geometric Approach

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For a given frequency, as **electron density increases** index of refraction **decreases**

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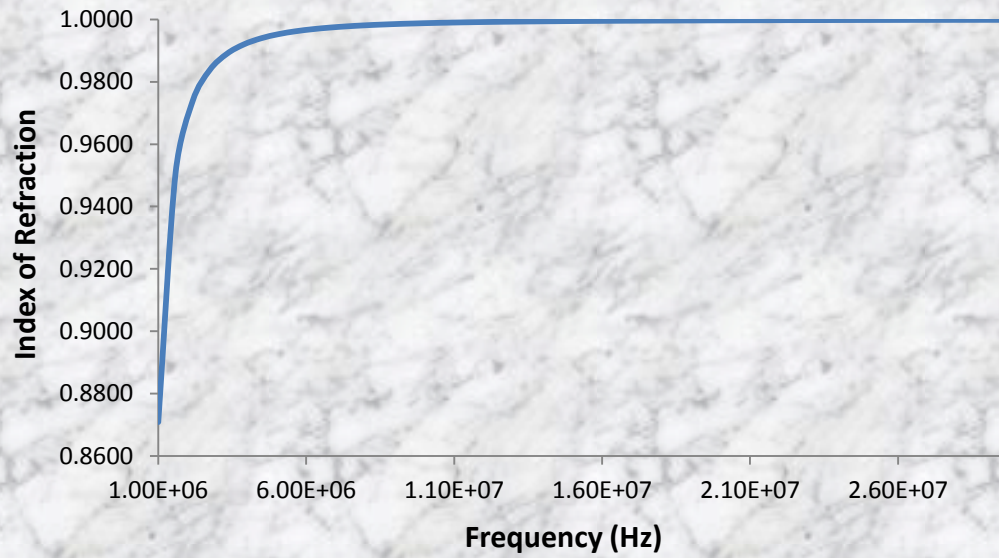
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Within a dispersive media such as the ionosphere:

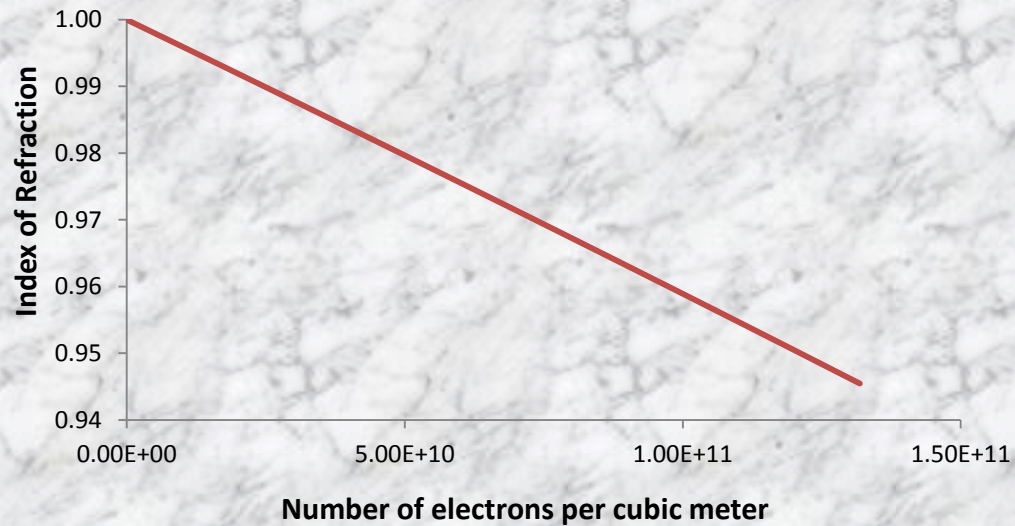
For a given frequency, as **electron density increases** index of refraction **decreases**

For a given electron density as **frequency increases** index of refraction approaches **unity**

Index of Refraction as Function of Frequency



Index of Refraction as Function of Electron Density



Geometric Approach

$$\mu = \sqrt{1 - \frac{\omega_p^2}{\omega^2}}$$

Three Cases:

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$$\omega \gg \omega_p \Rightarrow \mu \sim 1 \Rightarrow \text{Continues through ionosphere}$$

The Basic Absorption Equation

Basic Absorption Equation

Equation for total system loss

$$L_s = L_{ta} + L_{tp} - G_t + L_p + L_{rp} - G_r + L_{ra}$$

Note: Each term is a base-10 log. Hence, we add them

Basic Absorption Equation

Equation for total system loss

$$L_S = \overset{\text{Transmitting}}{\boxed{L_{ta} + L_{tp} - G_t}} + \boxed{L_p} + \overset{\text{Receiving}}{\boxed{L_{rp} - G_r + L_{ra}}}$$

What goes on in between

Critical term is path loss

Note: Each term is a base-10 log. Hence, we add them

Basic Absorption Equation

Equation for path loss

$$L_p = L_d + L_a + L_f$$

Basic Absorption Equation

Equation for path loss

$$L_p = \cancel{L_d} + L_a + \cancel{L_f}$$

Critical term is absorption.

Hence, we focus on the absorption term

Basic Absorption Equation

Equation for absorption

$$L_a = 10 \log \left(\frac{P_r}{P_u} \right)$$

$P_r \equiv$ Actual power received

$P_u \equiv$ Power received without absorption

Basic Absorption Equation

Equation for absorption

$$L_a = 20 \log(\rho)$$

$\rho \equiv$ Apparent Reflection Coefficient

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$\kappa \equiv$ Measure of amplitude decay per unit distance

$$\text{Let } \rho \equiv \frac{I}{I'_0} \Rightarrow \rho = \exp \left[- \int \kappa ds \right]$$

Basic Absorption Equation

$$\Rightarrow \ln(\rho) = - \int \kappa \, ds$$

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Since there are roughly 8.69 dB per neper the ***absorption equation*** has units of ***dB per unit length***

Types of Absorption

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Type of absorption depends on relationship between **radio wave frequency** and **plasma frequency**

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Deviative absorption **uncommon in D-region**

Types of Absorption

Type of absorption depends on relationship between **radio wave frequency** and **plasma frequency**

Type two: Radio wave frequency greater than plasma frequency

$$f > f_p$$

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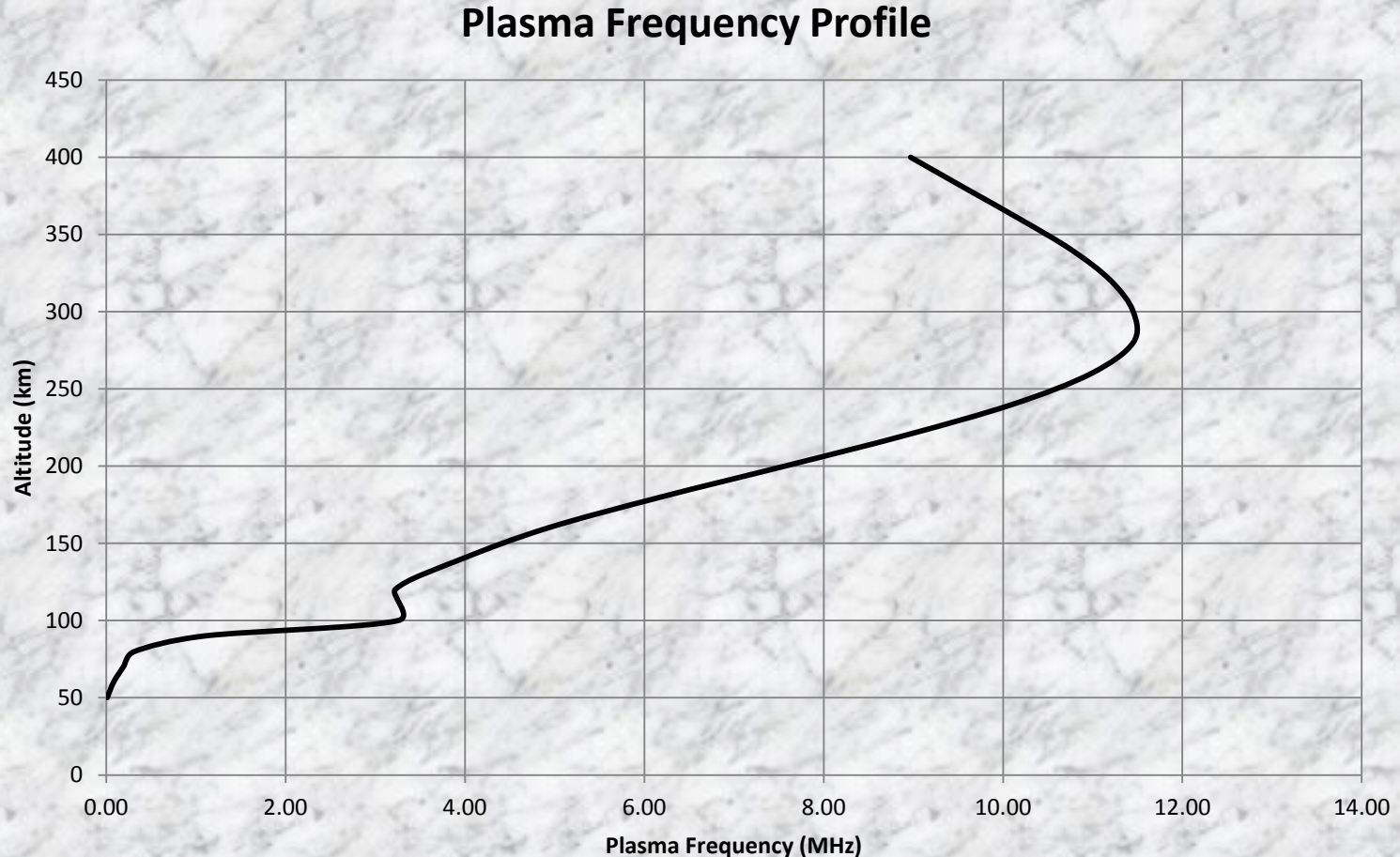
$$\Rightarrow \mu \approx 1$$

$$v_{phase} = \frac{c}{\mu} \approx c \text{ and } v_{group} = c\mu \approx c$$

Hence, wave propagates at about speed of light

This is called **non-deviative absorption**
Very common in D-region

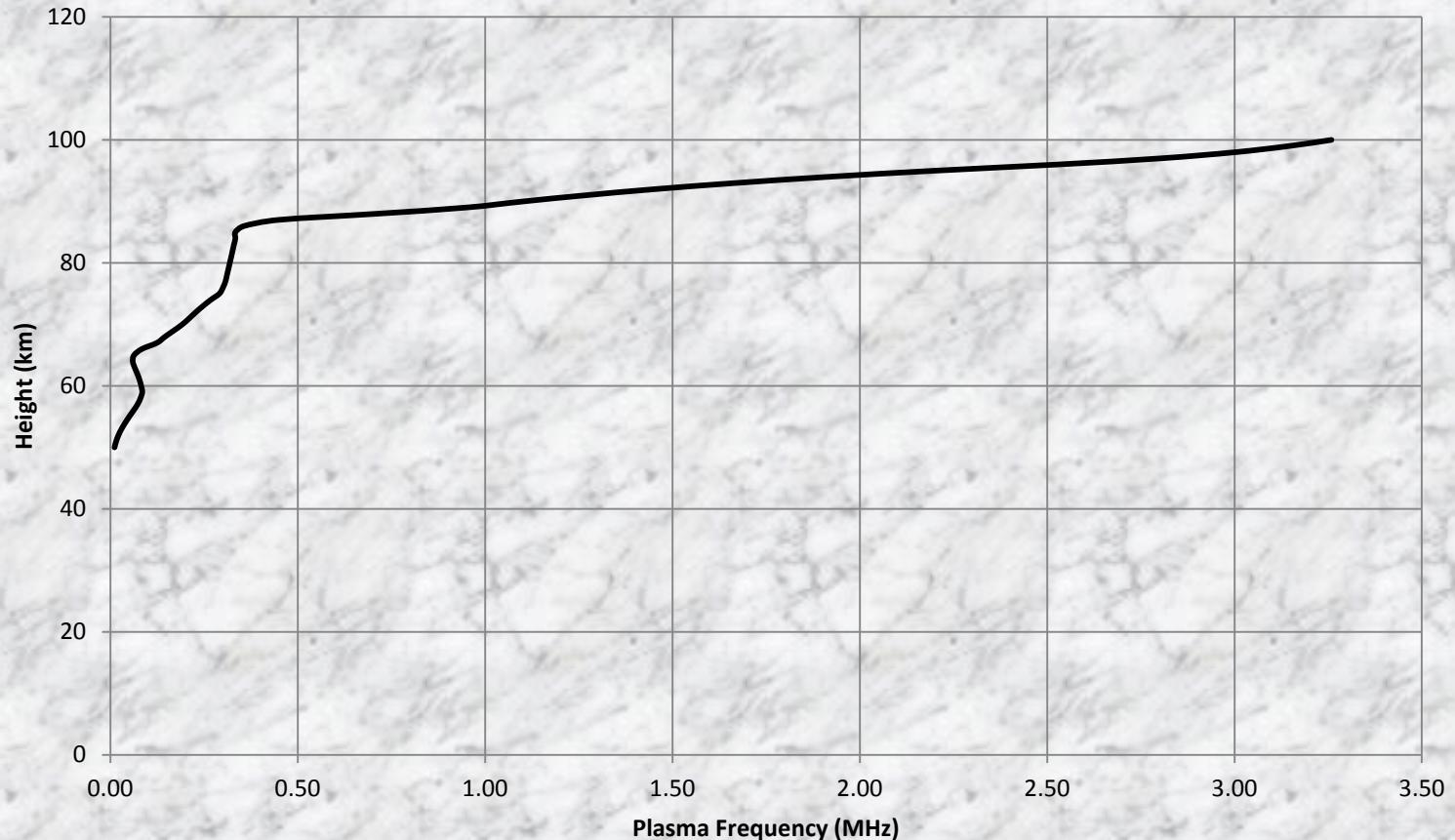
Types of Absorption



Profile of plasma frequency from 50-400 km. But we're really interested in D-region

Types of Absorption

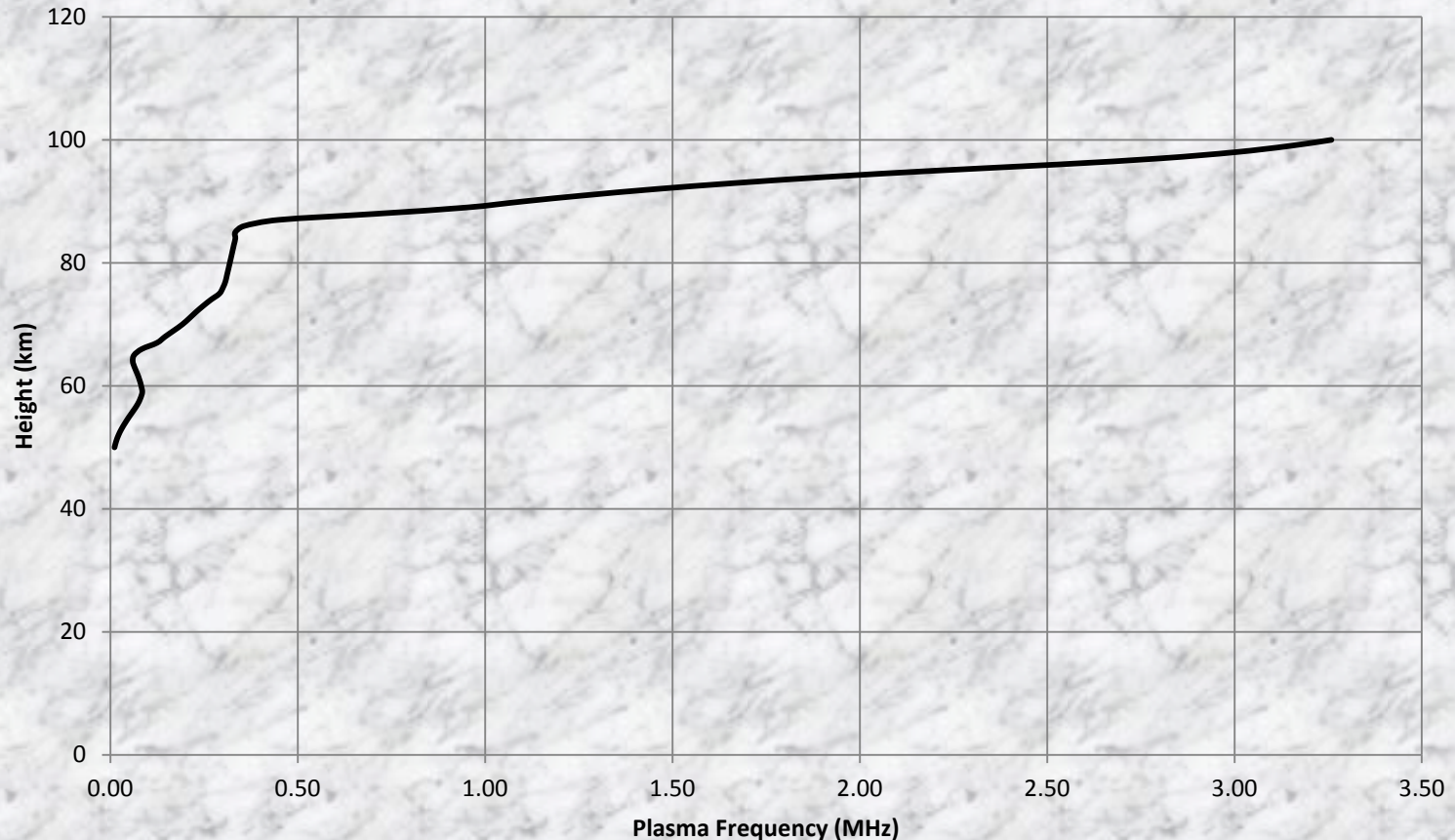
Typical Plasma Frequency Profile within D-region



Profile of plasma frequency from 50-100 km.
Note plasma frequency nearly always ***less than 3.0 MHz***

Types of Absorption

Typical Plasma Frequency Profile within D-region



Thus we see that non-deviative absorption dominates in the D-region
(Assuming $3.0 < f < 30.0 \text{ MHz}$)

Absorption Coefficient

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In the absorption equation *kappa* is defined as the ***absorption coefficient***

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Hence, the absorption equation is based on first principles**

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$$\kappa = \frac{e^2}{2\epsilon_0 mc} \left(\frac{1}{\mu} \right) \left(\frac{n_e \nu}{\omega^2 + \nu^2} \right)$$

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Units of *kappa* are ***nepers per unit length***

Defining Terms:

$e \equiv$ elementary charge = $1.6 \times 10^{-19} \text{ C}$

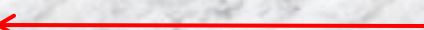
$\epsilon_0 \equiv$ permittivity of free space = $8.85 \times 10^{-12} \text{ F} \cdot \text{m}^{-1}$

$m_e \equiv$ electron mass = $9.11 \times 10^{-31} \text{ kg}$

$c \equiv$ speed of light in vacuum = $3.0 \times 10^8 \text{ m} \cdot \text{s}^{-1}$

$\mu \equiv$ real part of refractive index $\cong 1.0$

$n_e \equiv$ electron density

$\nu \equiv$ collision frequency  **New important term!**

$\omega \equiv$ angular frequency of wave

Absorption Coefficient

$$\kappa = \frac{e^2}{2\epsilon_0 mc} \left(\frac{1}{\mu} \right) \left(\frac{n_e \nu}{\omega^2 + \nu^2} \right)$$

Units of *kappa* are ***nepers per unit length***

Plugging in constant values we find that,

$$\kappa \cong 5.29 \times 10^{-6} \left(\frac{n_e \nu}{\omega^2 + \nu^2} \right) \text{ Np} \cdot \text{m}^{-1}$$

Absorption Equation Revisited

Absorption Equation

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$$\Rightarrow L_a = -4.60 \times 10^{-5} \int \left(\frac{n_e \nu}{\omega^2 + \nu^2} \right) ds \text{ dB} \cdot \text{m}^{-1}$$

Absorption Equation

$$L_a = -4.60 \times 10^{-5} \int \left(\frac{n_e \nu}{\omega^2 + \nu^2} \right) ds$$

In this form integral is over ***path length***

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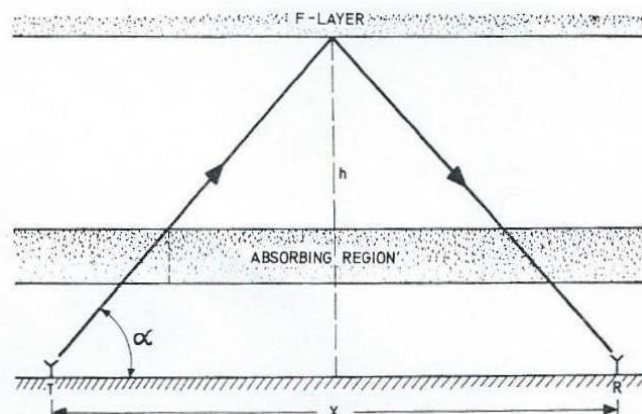
Electron density and collision frequency can be functions of ***height***

Absorption Equation

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Electron density and collision frequency can be functions of **height**



$$\sin \alpha = \frac{\text{opposite}}{\text{hypotenuse}} = \frac{h}{s}$$

Absorption Equation

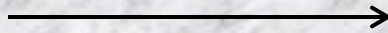
$$L_a = -4.60 \times 10^{-5} \int \left(\frac{n_e v}{\omega^2 + v^2} \right) ds$$

$$\Rightarrow s = \frac{h}{\sin \alpha}$$

Absorption Equation

$$L_a = -4.60 \times 10^{-5} \int \left(\frac{n_e v}{\omega^2 + v^2} \right) ds$$

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Now integrated over height

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For special case of vertical transmission:

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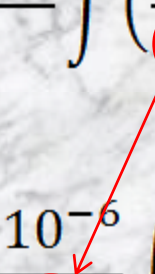
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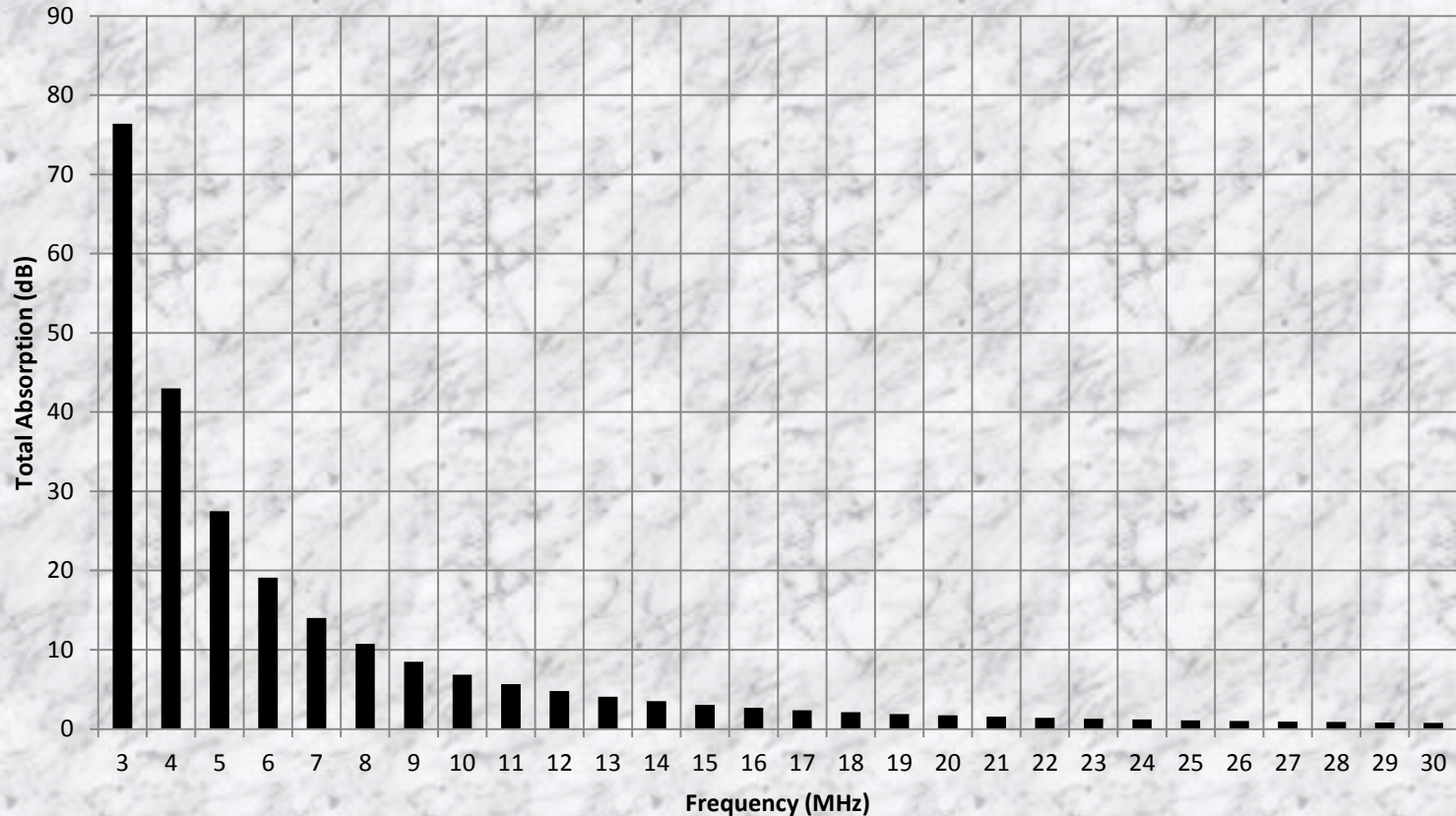
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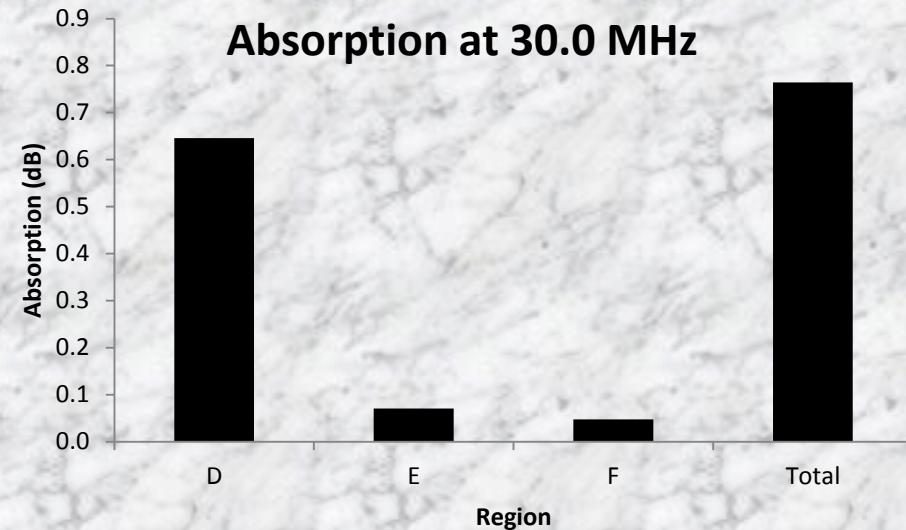
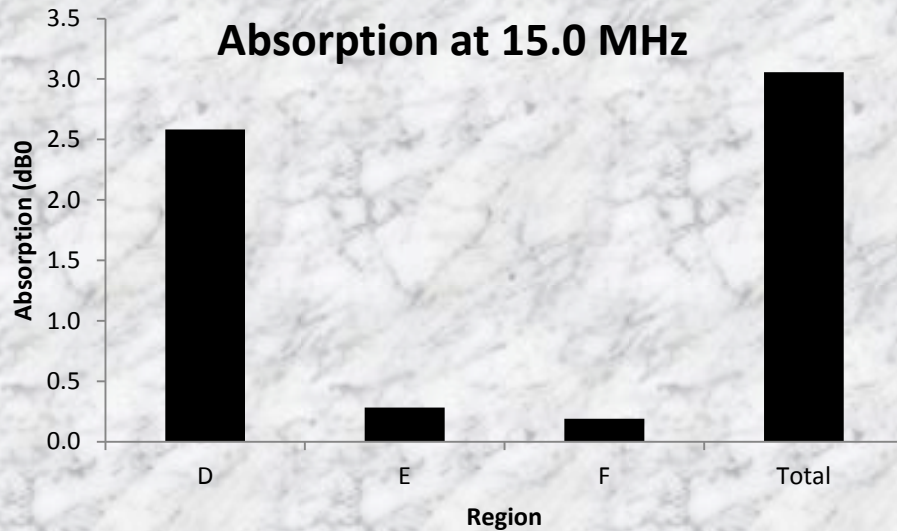
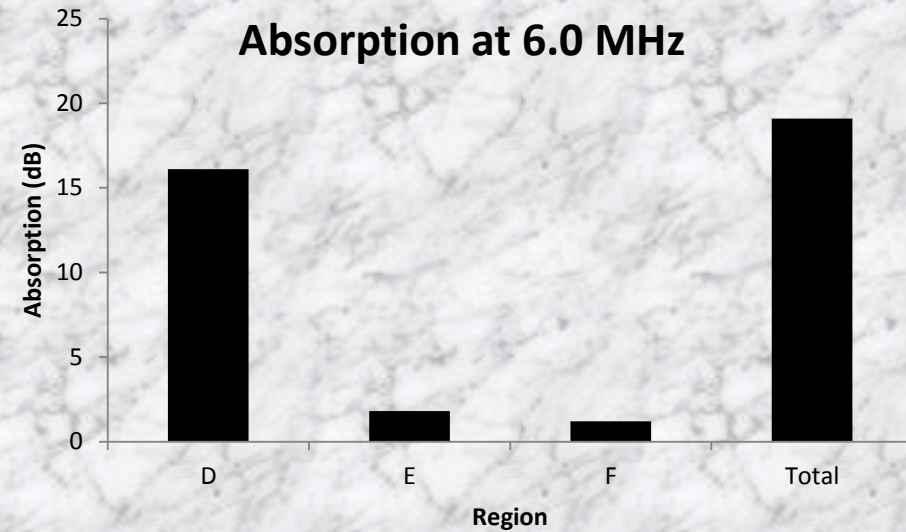
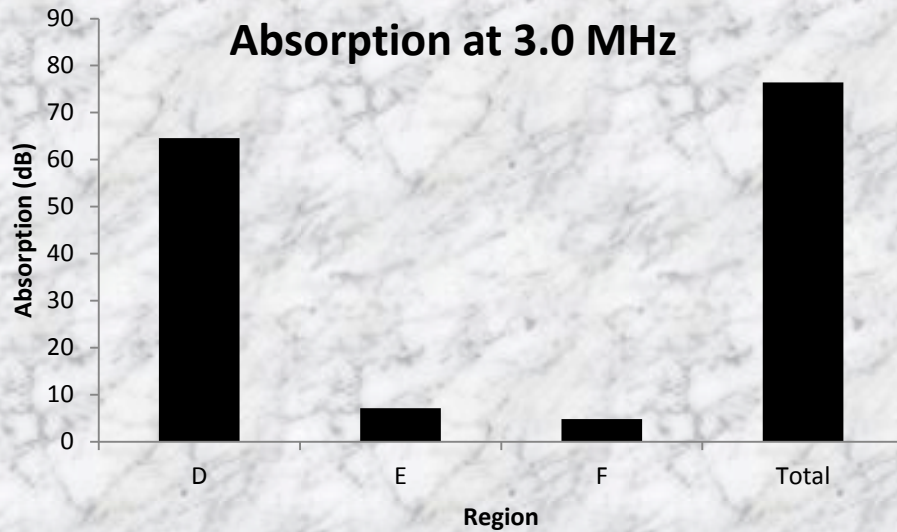
The absorption equation in terms of radio wave frequency in **cycles per second**.

Absorption Equation

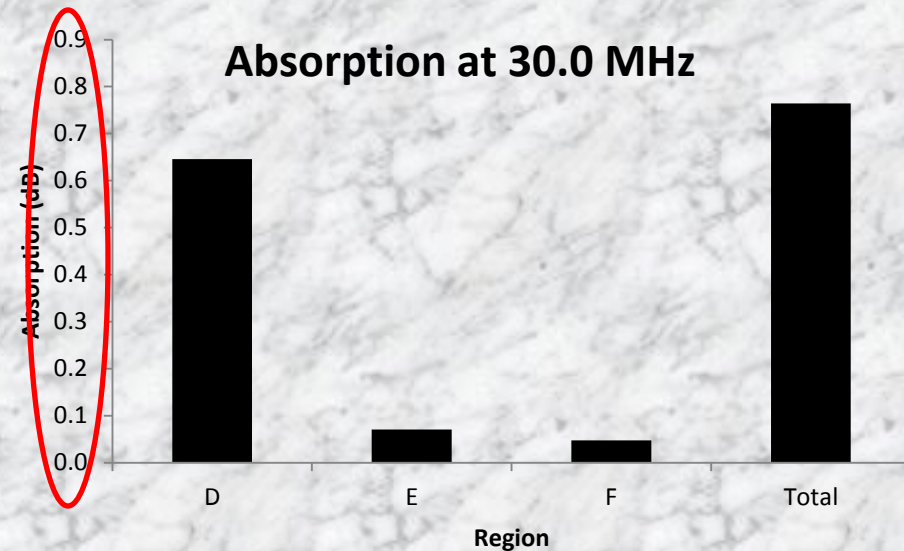
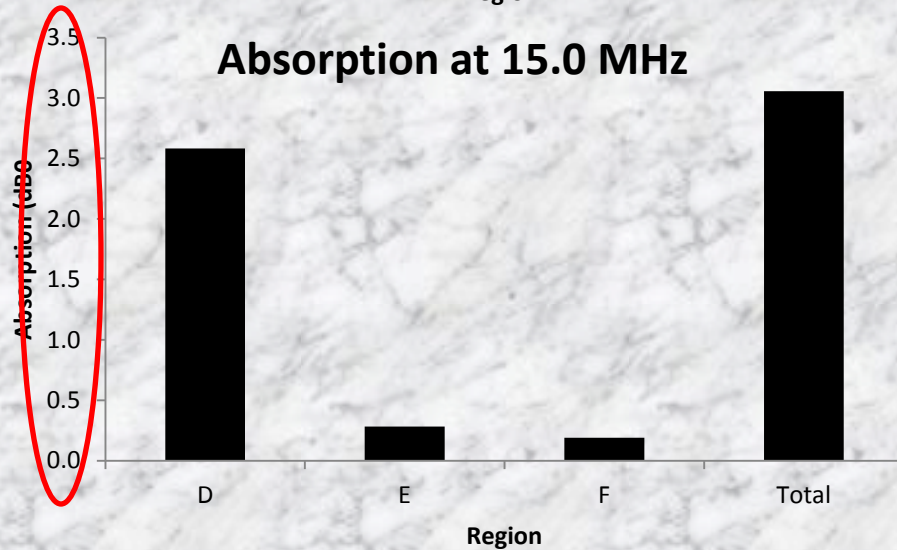
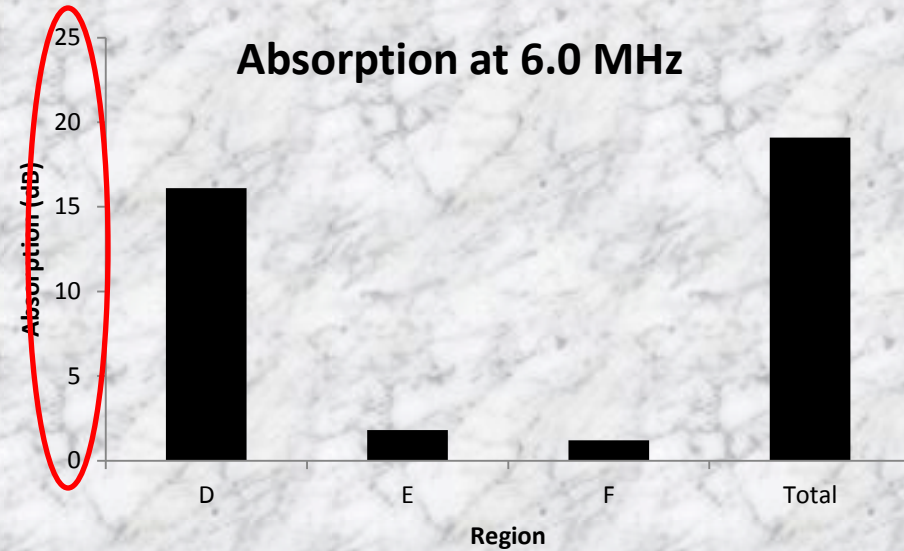
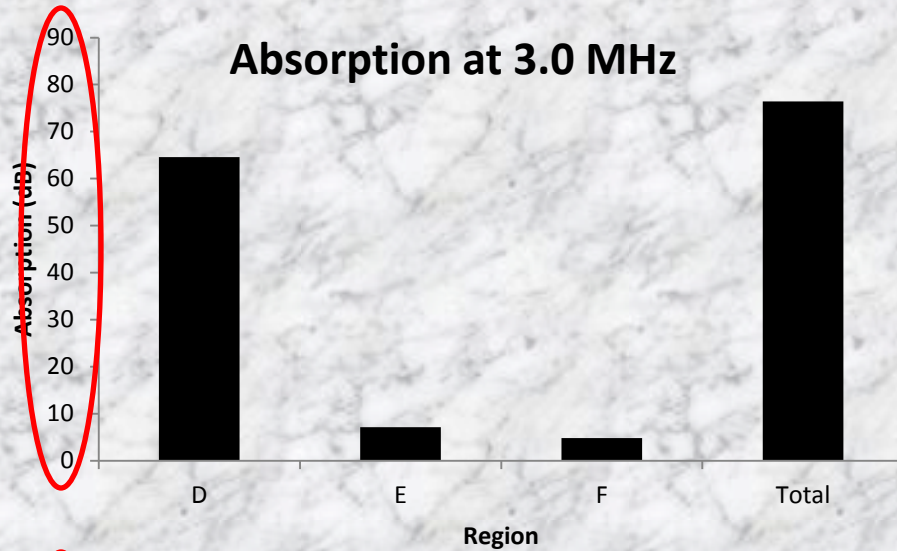
Total Absorption by Frequency



Thus we see the **frequency-dependence** of absorption



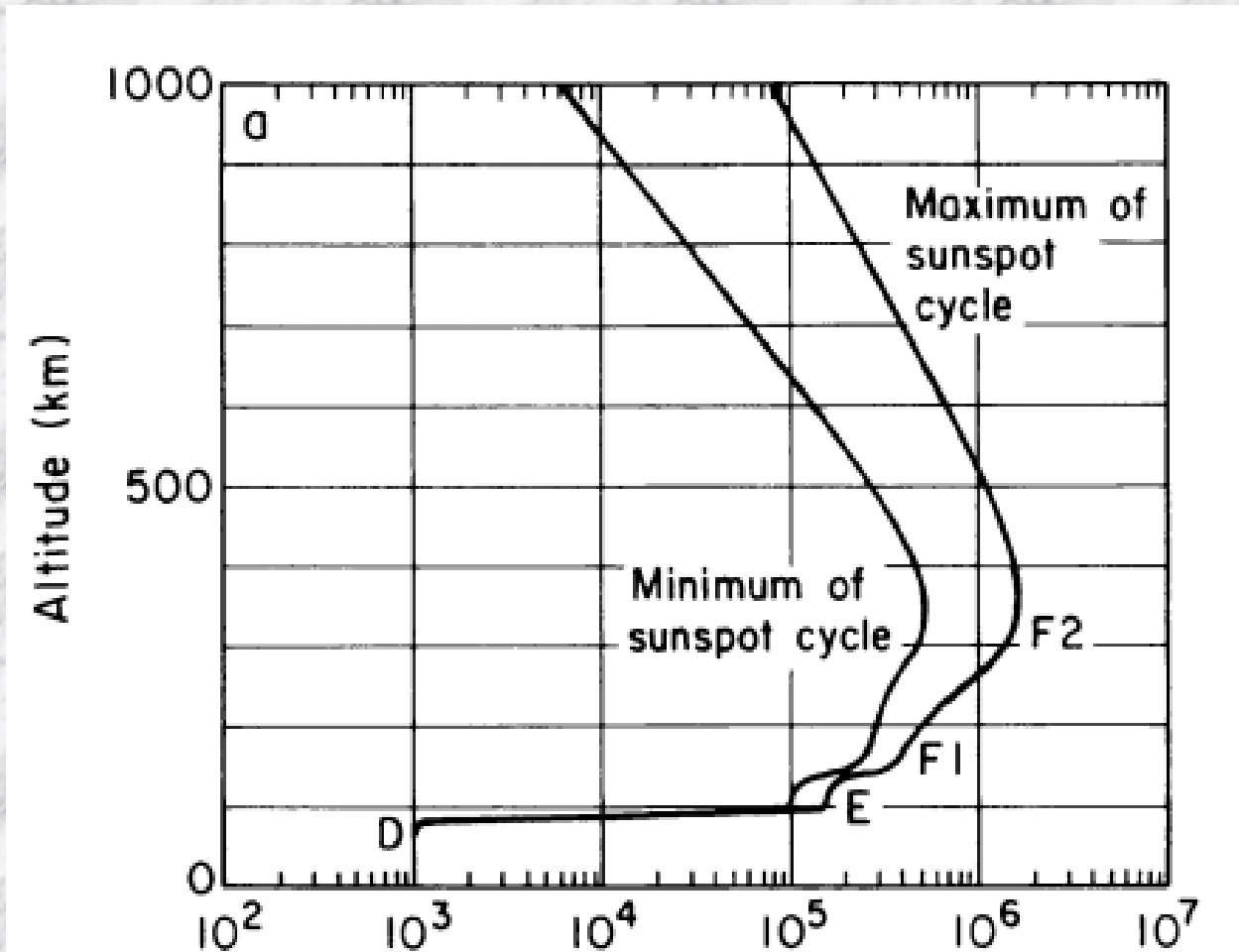
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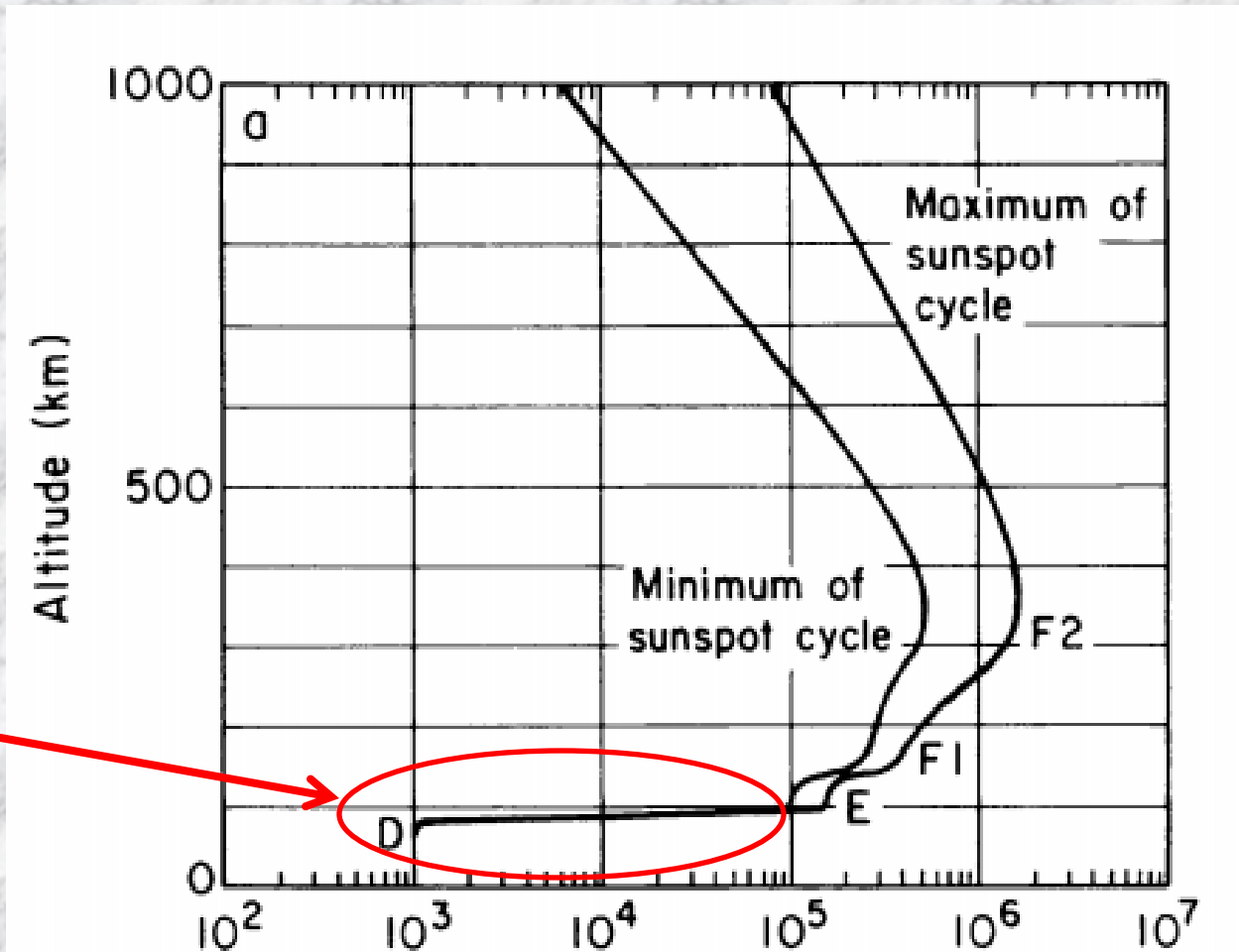
Electron Density

Electron Density



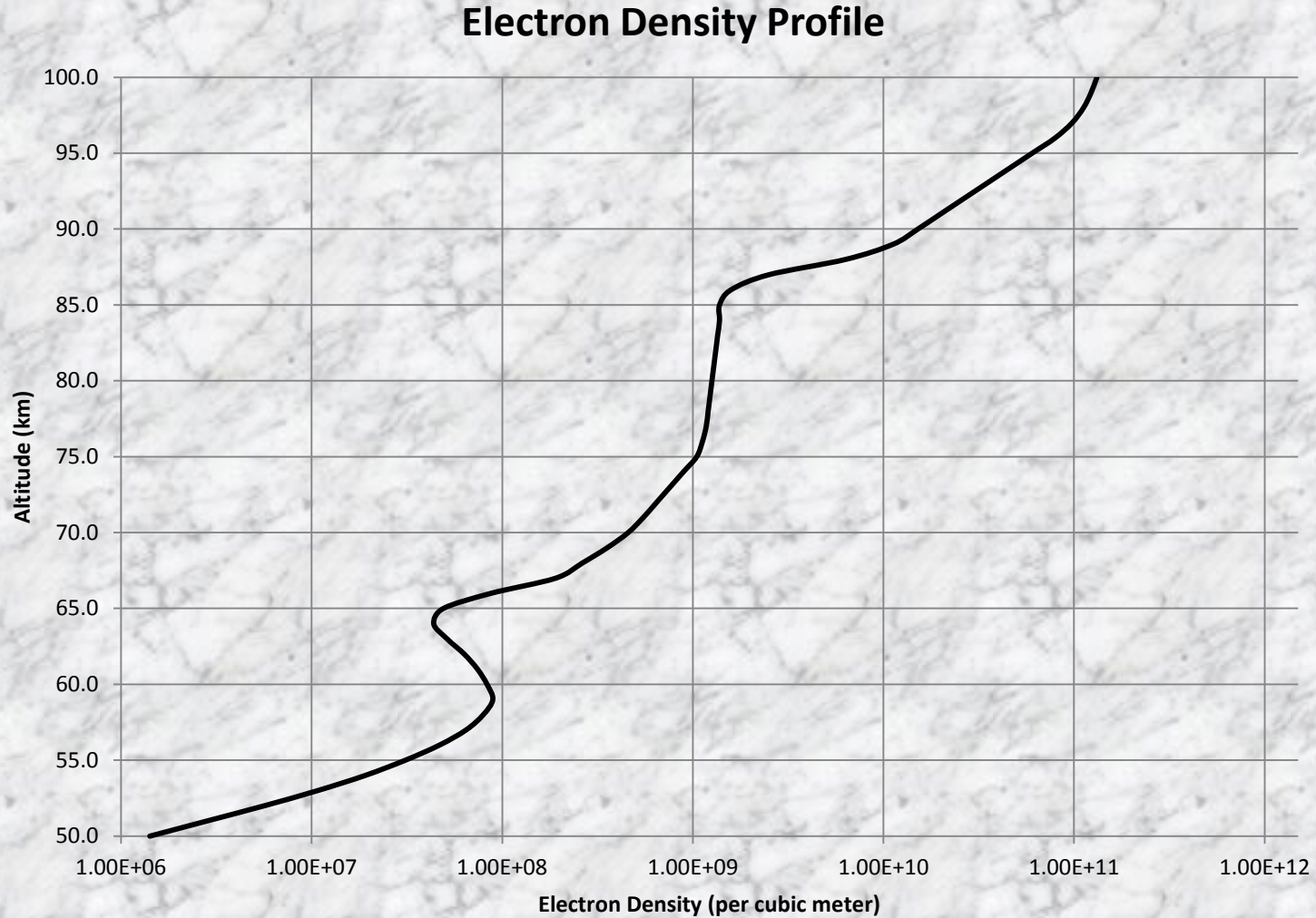
Typical electron concentration per cubic centimeter (Daytime)

Electron Density



Typical electron concentration per cubic centimeter (Daytime)

Electron Density



Electron density profile below 100 km

Electron Density

Everything depends on electron density

Electron Density

Everything depends on electron density

- **Plasma Frequency**

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Thus we see that electron density is the **most critical** component

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In the above T_e is the electron temperature

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- Most HF absorption takes place within the D-region
- Within the D-region non-deviative absorption dominates
- **The electron density is the most critical component**

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We showed the following to be true:

Non-deviative absorption within the D-region can be described mathematically in terms of neutral density or collision frequency,

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Acknowledgments

Special thanks to the following who assisted in the preparation of the presentation

Dr. Jan J. Sojka

Dr. Vince Eccles

And thank you to my supervisory committee:

Doctors J. Sojka, D. Peak, B. Fejer, M. Taylor, R. Fullmer

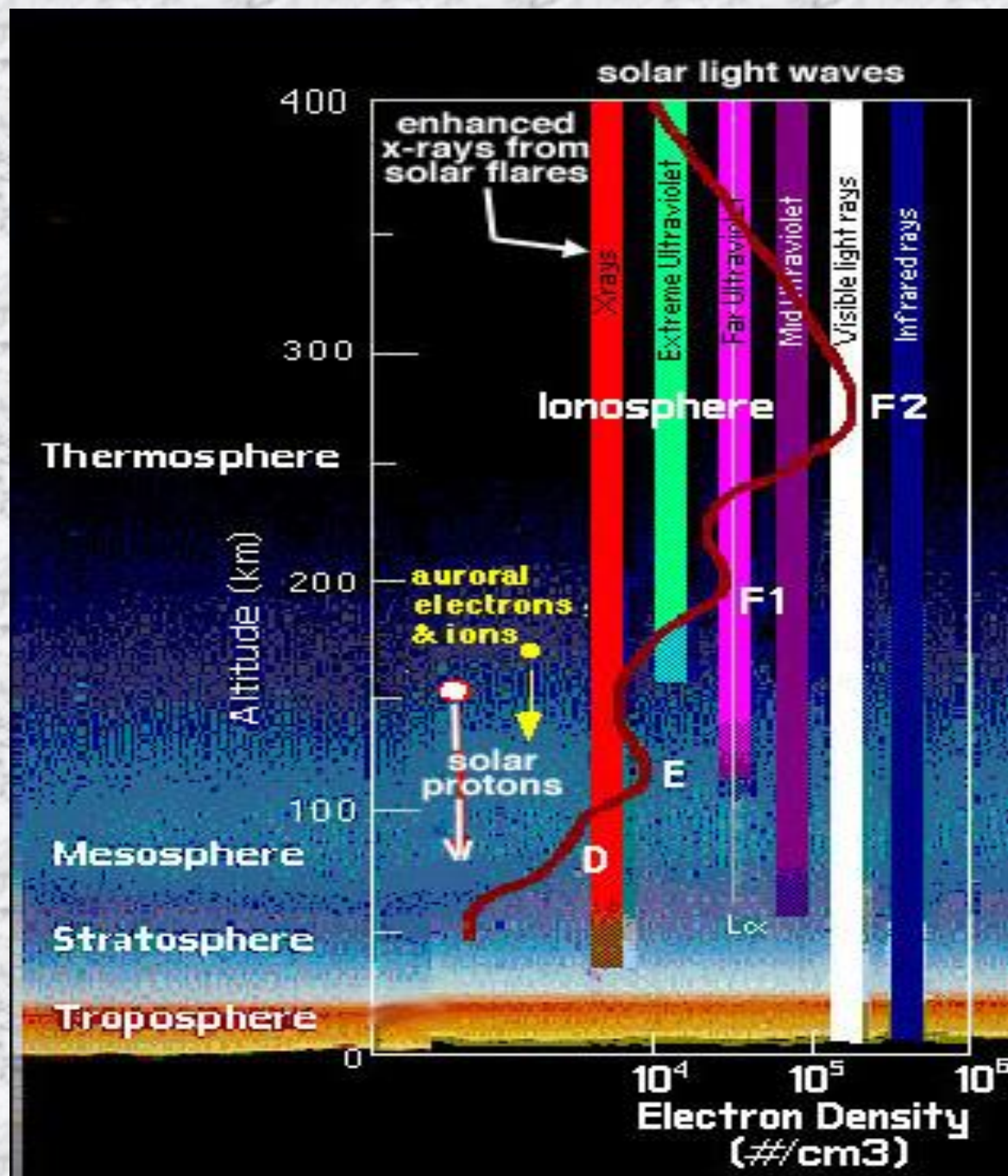
References

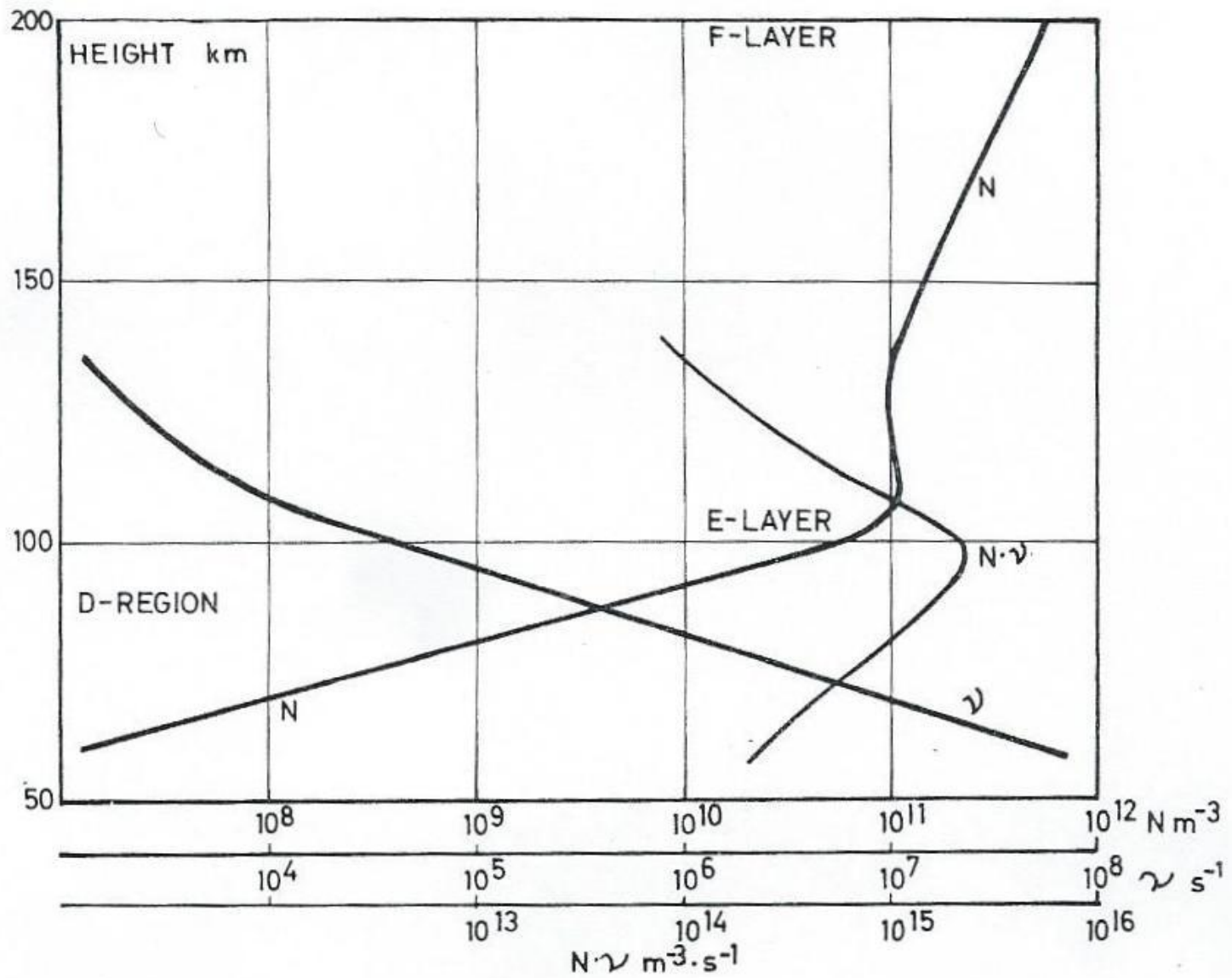
Key Sources

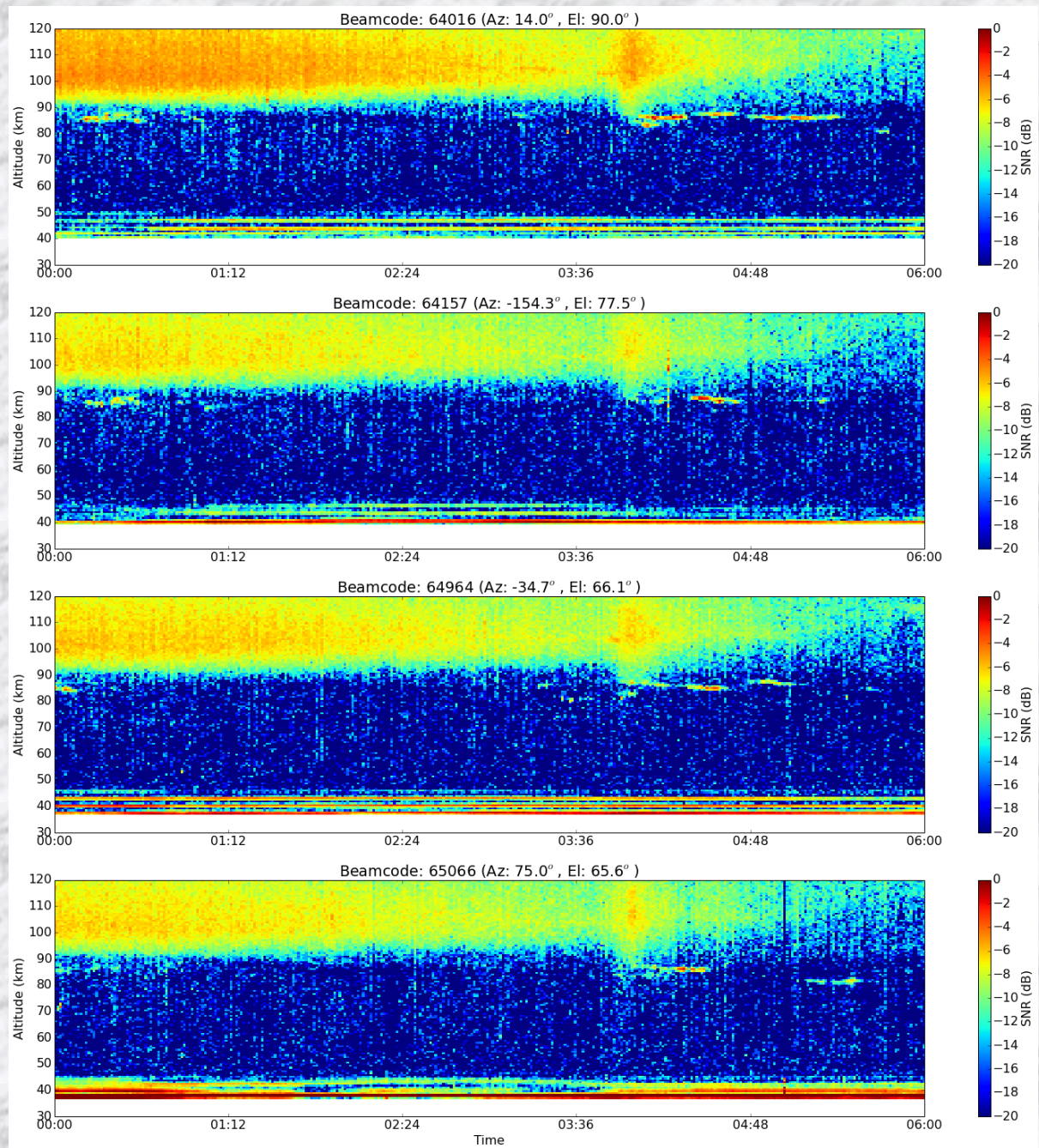
- ARRL, 1991. Chapter 23, Radio Wave Propagation. In: *The ARRL Antenna Book*. 16 ed. Newington: The American Radio Relay League, pp. 23.1-23.27.
- Bain, W. C. & Harrison, M. D., 1972. Model ionosphere for D region at Summer noon during sunspot maximum. *Proc. IEEE*, 119(7), pp. 790-796.
- Davies, K., 1965. Chapter 2: Theory of Wave Propagation. In: *Ionospheric Radio Propagation*. s.l.:National Bureau of Standards, pp. 45-100.
- Davies, K., 1965. Chapter 3, Synoptic Studies of the Ionosphere. In: *Ionospheric Radio Propagation*. s.l.:National Bureau of Standards, pp. 101-158.
- Davies, K., 1965. Chapter 5, Signal Strength. In: *Ionospheric Radio Propagation*. s.l.:National Bureau of Standards, pp. 217-256.
- Griffiths, D. J., 2013. *Introduction to Electrodynamics*. 4th ed. Boston: Pearson.
- Hunsucker, R. & Hargreaves, J., 2003. *The High-Latitude Ionosphere and its Effects on Radio Propagation*. First ed. Cambridge: Cambridge University Press.
- Jackson, J. D., 1999. *Classical Electrodynamics*. 3rd ed. Danvers: John Wiley & Sons, Inc.
- Kelley, M. C. & Heelis, R. A., 1989. Appendix B; Reference Material and Equations. In: *The Earth's Ionosphere: Plasma Physics and Electrodynamics*. San Diego: Academic Press, Inc, pp. 459-471.
- Lied, F. (., 1967. *High Frequency Radio Communications with Emphases on Polar Problems*. s.l.:The Advisory Group for Aerospace Research and Development.
- Moen, J., 2004. Chapter 3: Structure and Composition of the Middle and Upper Atmosphere. In: *FYS3610 Fall 2004 Space Physics Lecture Series*. Oslo: s.n.
- Moen, J., 2004. Chapter 4, The Ionosphere. In: *FYS3620 Fall 2004 Space Physics Lecture Series*. Oslo, Norway: s.n.
- Odenwald, S., 2010. Introduction to Space Storms and Radiation. In: C. J. Schrijver & G. L. Siscoe, eds. *Heliophysics, Space Storms and Radiation: Causes and Effects*. New York: Cambridge University Press, pp. 15-41.
- Rawer, K., 1952. Calculation of Sky-wave Field Strength. *The Wireless Engineer*, Volume 19, pp. 287-301.
- Schunk, R. & Nagy, A., 2009. *Ionospheres: Physics, Plasma Physics, and Chemistry*. Second ed. Cambridge: Cambridge University Press.
- Zolesi, B. & Cander, L., 2014. Chapter 2: The General Structure of the Ionosphere. In: *Ionospheric Prediction and Forecasting*. s.l.:Springer, pp. 11-47.

fin

Questions?







Frequency (MHz)	L_{α} (dB)
3	56.6
5	20.4
7	10.4
9	6.29
11	4.21
13	3.02
15	2.27
17	1.76
19	1.41
21	1.16
23	0.964
25	0.816
27	0.699
29	0.606
31	0.530

D-region absorption values using data from *Bain and Harrison*