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September 20, 2008

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TIME-RESOLVED AU-SCALE JETS TRACED BY MASERS IN THE IRAS 4A/B REGIONS OF NGC 1333

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Received 2007 August 30; accepted 2008 May 27

ABSTRACT

We present results of VLBI observations of the water masers associated with IRAS 4A and IRAS 4B in the NGC 1333 star-forming region taken in four epochs over a two month period. Both objects have been classified as extremely young sources, and each source is known to be a multiple system. Using the Very Long Baseline Array, we detected 35 masers in epoch I, 40 masers in epoch II, 35 in epoch III, and 24 in epoch IV. Only one identified source in each system is associated with these masers. These data are used to calculate proper motions for the masers and trace the jet outflows within 100 AU of IRAS 4A2 and IRAS 4BW. In IRAS 4A2 there are two groups of masers, one near the systemic cloud velocity and one redshifted. They expand linearly away from each other at velocities of 53 km s^{-1} . In IRAS 4BW, masers are observed in two groups that are blueshifted and redshifted relative to the cloud velocity. They form complex linear structures with a thickness of 3 mas (1 AU at a distance of 320 pc) that expand linearly away from each other at velocities of 78 km s^{-1} . Neither of the jet outflows traced by the maser groups align with the larger scale outflows. We suggest the presence of unresolved companions to both IRAS 4A2 and 4BW.

Subject headings: ISM: jets and outflows — masers — shock waves — stars: mass loss — stars: pre-main-sequence

1. INTRODUCTION

Water masers are excellent probes of astrophysical flows. Their large flux densities and exceedingly compact sizes make them perfect Very Long Baseline Interferometer (VLBI) targets. Many masers observed in star-forming regions are thought to form in shocks either in or along the outflow of material commonly seen associated with young stellar objects (YSOs). They provide a high-resolution probe of the base of the stellar wind that is unaffected by extinction from dust or extensive interaction with ambient material. The milliarcsecond resolution of VLBI is critical to resolve the individual maser components and outflows, particularly in high stellar density environments.

Long-term total power monitoring has shown that masers around low-mass YSOs are more episodic than those around higher mass YSOs; however, their detectable phases are sufficiently long for proper-motion measurements if the observations are spaced by no more than 2–3 weeks (Furuya et al. 2003; Brand et al. 2003; Claussen et al. 1996; Wilking et al. 1994). Furthermore, if they arise in the warm (400 K) spatially confined postshock region, as postulated in various models (Elitzur et al. 1989, 1992; MacLow et al. 1994; MacLow & Elitzur 1992), they should have space velocities sufficient to produce measurable proper motions over 1–3 weeks. However, regular monitoring of low-mass YSOs is necessary to ensure that VLBI observations are conducted during periods of high maser activity.

Using the NRAO¹ Very Long Baseline Array (VLBA), proper-motion studies have been made of water masers associated with high-mass and intermediate-mass YSOs such as Cepheus A HW2 (Torrelles et al. 2001), S106 FIR (Furuya et al. 2000), NGC 2071 (Seth et al. 2002), and IRAS 20050+2720 MMS1 (Furuya et al. 2005), as well as with lower mass YSOs such as Serpens SMM1, RNO 15-FIR, and IRAS 05413–0104 (Moscadelli et al. 2006; Claussen et al. 1998). For example, water masers in IRAS 05413–0104 (associated with HH 212 at a distance of 450 pc), were found to lie in a structure of 10 AU in length and less than 0.5 AU thick. The masers appeared to arise in shock-related structures which showed proper motions along the axis of the outflow of 60 km s^{-1} and displayed coherent structures over timescales of 2–3 weeks (but which was less discernible over timescales as long as several months). Comparison with molecular and infrared observations of IRAS 05413–0104 clearly demonstrated that proper-motion studies of water masers in the region (within 40 AU) of the central source are among the best tools for studying the kinematics of the jets emanating from embedded YSOs.

IRAS 4, comprising multiple sources whose submillimeter-dominated “Class 0” spectral energy distributions suggest extreme

¹ The National Radio Astronomy Observatory is a facility of the National Science Foundation operated under cooperative agreement by Associated Universities, Inc.

youth, is located in the NGC 1333 star-forming region. As a nearby active and multiple young object, it presents a prime candidate for VLBI observations of its associated water masers. Located at a distance of 320 pc (de Zeeuw et al. 1999), the NGC 1333 molecular cloud hosts a double infrared cluster of about 200 YSOs identified in near-infrared surveys (Strom et al. 1976; Aspin et al. 1994; Lada et al. 1996; Wilking et al. 2004) and X-ray surveys (Preibisch 1997, 2003; Getman et al. 2002). Far-infrared observations have revealed not only the higher luminosity sources in NGC 1333, but also YSOs in the earliest phase of evolution, such as the Class 0 sources IRAS 4 and IRAS 2 (Harvey et al. 1984; Jennings et al. 1987). IRAS 4 was found to be a binary at submillimeter wavelengths (Sandell et al. 1991); hence IRAS 4A and 4B. IRAS 4A is located $31''$ to the northwest of 4B at a position angle of -45° . Subsequently, both IRAS 4A and 4B were found to be binary systems. Lay et al. (1995) performed single-baseline interferometry in the submillimeter and found that 4A was a binary with an angular separation of $1.8''$ (or 580 AU at 320 pc). Interferometric observations in the millimeter and radio continuum confirmed the separation and determined that the position angle of separation was $\approx -20^\circ$ (Looney et al. 2000; Reipurth et al. 2002). While the easternmost source (IRAS 4A1) dominates the millimeter and radio continuum emission, the westernmost source (IRAS 4A2) appears to be relatively more evolved, with a warm ammonia core (Wootten & Mangum 1993; Shah 2000) and a wealth of complex organic molecules (Bottinelli 2007, private communication). Submillimeter and millimeter wave continuum observations have confirmed the prediction of Lay et al. (1995) that IRAS 4B is also a binary with components (4BE and 4BW) separated by about $10''$ (3200 AU in projection) along an east-west direction (Smith et al. 2000; Looney et al. 2000; Sandell & Knee 2001).² As in the case of IRAS 4A2, complex organic molecules have been detected toward IRAS 4B (Bottinelli et al. 2007), presumably from a hot core associated with 4BW. We note that IRAS 4BE has yet to be detected in the radio continuum (Reipurth et al. 2002) or demonstrate any sign of outflow activity. Infall motions have been detected toward IRAS 4A and 4B (Di Francesco et al. 2001), unambiguously toward the former, and strongly indicated toward the latter.

The detection of water masers toward IRAS 4A and 4B suggests that they are YSOs also associated with mass outflow. Water masers were first detected in single-dish observations by Haschick et al. (1980) and later monitored by Claussen et al. (1996). These observations established the maser emission as highly variable on monthly timescales, sometimes being completely absent, while at other times reaching peaks of 10 Jy. Detected emission was found from -10 to $+15$ km s $^{-1}$, but the angular resolution was not sufficient to determine which masers were associated with 4A and 4B. Beginning in 1983, VLA observations have shown water maser activity associated only with IRAS 4A2 and/or 4BW (Rodríguez et al. 2002; H. A. Wootten 1998, unpublished; Furuya et al. 2003). VLBA observations of the water masers in IRAS 4A and 4B, acquired in 2003, have been presented by Desmurs et al. (2006). While masers were detected in only two of the four epochs observed, they confirmed the association of masers with IRAS 4A2, and expansion motion between the two spots was detected. In IRAS 4BW, the five maser spots detected were redshifted and formed a chain 80 milliarcsecond (mas) in extent, with some masers displaying proper motion toward the north.

Molecular outflows have been mapped toward both IRAS 4A and 4B. IRAS 4A is associated with a highly collimated outflow about 20,000 AU in extent seen in HCN and SiO, with a position angle close to 20° (Girart et al. 1999; Choi 2001, 2005). The origin of the outflow is likely IRAS 4A2 (Choi 2005). On a larger scale ($\approx 1 \times 10^5$ AU), the outflow defined by CO and molecular hydrogen has a position angle of 45° (Blake et al. 1995; Choi et al. 2006). The shift in position angle from small to large scale is perhaps due to a combination of a shift in the direction of the magnetic field, an encounter with denser ambient gas, and precession of the outflow axis (Choi et al. 2006). A more compact outflow is associated with IRAS 4BW, with a position angle close to 0° (Blake et al. 1995; Choi 2001) with perhaps a second outflow with a position angle of -35° (Di Francesco et al. 2001).

In this paper we present VLBA observations of water masers associated with the IRAS 4 region obtained over four epochs in 1998 and spaced by about 1 month. We also describe Very Large Array (VLA) observations, taken 1 month in advance of the start of the VLBA observations, that showed IRAS 4 to be a very active maser source and yielded absolute positions for the masers and the 1.3 cm continuum sources IRAS 4A1 and IRAS 4BW. With the VLBA, we detected 35 masers in epoch I, 40 masers in epoch II, 35 in epoch III, and 24 in epoch IV associated with IRAS 4A2 and IRAS 4BW. We use these data to reveal the structure of shock fronts associated with stellar winds from these YSOs and to estimate proper motions for masers detected in all four epochs. The origin of the maser emission and its relationship to the stellar winds and larger scale molecular outflows from these YSOs is discussed.

2. OBSERVATIONS AND REDUCTIONS

2.1. VLA Observations

The VLA was used on 1998 August 12 to observe the water masers toward IRAS 4A and IRAS 4B to check that the masers were strong enough to observe subsequently with the VLBA. The array was in the B configuration. The phase-tracking center for the array was set to be midway between the IRAS 4A and IRAS 4B positions, spread over a range of hour angles to optimize the u, v coverage. Two observation modes were used. In the first, observations were made with both circular polarizations, a total bandwidth of 1.5625 MHz (total velocity extent of 21 km s $^{-1}$), and a correlator setup that gave 127 spectral channels, each with width of 12.2 kHz (0.16 km s $^{-1}$). In this mode the water masers were observed with reasonably high spectral resolution in order to spectrally resolve the masers, and the observations lasted for ~ 42 minutes (on-source).

In the second mode, the correlator was configured for continuum observations with one polarization pair (of IFs) set to a center frequency of 22264.9 MHz and a bandwidth of 50 MHz, while the other polarization pair was set to a bandwidth of 0.781 kHz, and the center frequency adjusted so that the strongest masers were placed at the center of the narrow band. This setup allowed the phases to be tracked using the strong signal in the narrow band, in order to calibrate tropospheric phase changes to apply to the 50 MHz continuum band (away from the maser emission). In this mode the continuum emission from the IRAS 4 region was observed for ~ 3.5 hr (on-source).

The data were edited, calibrated, and imaged with the NRAO Astronomical Image Processing System (AIPS; Greisen 2003) in the standard manner. In the high spectral resolution data, Stokes I images were made using uniform weighting of the u, v data, of each spectral channel, and were CLEANed to an rms noise (in channels without strong maser features) of ~ 20 mJy beam $^{-1}$.

² We adopt the naming convention introduced by Sandell & Knee (2001). The western (eastern) source is referred to as IRAS 4BI (IRAS 4BII) by Smith et al. (2000) as IRAS 4B (IRAS 4C) by Looney et al. (2000) and as IRAS 4B (4B') by (Di Francesco et al. 2001).

TABLE 1
IRAS 4A AND IRAS 4B PROPERTIES

Property	Value	Reference
IRAS 4A		
A1 2.7 mm position (J2000.0).....	$3^{\text{h}}29^{\text{m}}10.510^{\text{s}}, +31^{\circ}13'31.01''$	Looney et al. (2000)
A1 1.3 cm position (J2000.0).....	$3^{\text{h}}29^{\text{m}}10.527^{\text{s}}, +31^{\circ}13'31.06''$	This study
A1 3.6 cm position (J2000.0).....	$3^{\text{h}}29^{\text{m}}10.529^{\text{s}}, +31^{\circ}13'31.05''$	Reipurth et al. (2002)
A2 2.7 mm position (J2000.0).....	$3^{\text{h}}29^{\text{m}}10.413^{\text{s}}, +31^{\circ}13'32.20''$	Looney et al. (2000)
A2 3.6 cm position (J2000.0).....	$3^{\text{h}}29^{\text{m}}10.421^{\text{s}}, +31^{\circ}13'32.21''$	Reipurth et al. (2002)
A2 H ₂ O NW position (J2000.0).....	$3^{\text{h}}29^{\text{m}}10.4101^{\text{s}}, +31^{\circ}13'32.235''$	VLA, 1998 Aug 12; this study
A2 H ₂ O SE position (J2000.0).....	$3^{\text{h}}29^{\text{m}}10.4226^{\text{s}}, +31^{\circ}13'32.142''$	VLA, 1998 Aug 12; this study
A1 peak flux (2.7 mm).....	107 mJy	(2'') Looney et al. (2000)
A2 peak flux (2.7 mm).....	23 mJy	(2'') Looney et al. (2000)
A1 peak flux (3.6 cm).....	0.32 mJy	Reipurth et al. (2002)
A2 peak flux (3.6 cm).....	0.11 mJy	Reipurth et al. (2002)
Bolometric luminosity	$14 L_{\odot}$	André et al. (2000)
Bolometric temperature.....	34 K	André et al. (2000)
IRAS 4B		
2.7 mm position (J2000.0).....	$3^{\text{h}}29^{\text{m}}11.988^{\text{s}}, +31^{\circ}13'08.10''$	Looney et al. (2000)
1.3 cm position (J2000.0).....	$3^{\text{h}}29^{\text{m}}12.001^{\text{s}}, +31^{\circ}13'08.16''$	This study
3.6 cm position (J2000.0).....	$3^{\text{h}}29^{\text{m}}12^{\text{s}}.003^{\text{s}}, +31^{\circ}13'08.14''$	Reipurth et al. (2002)
B H ₂ O NW position (J2000.0).....	$3^{\text{h}}29^{\text{m}}11.9923^{\text{s}}, +31^{\circ}13'08.363''$	VLA, 1998 Aug 12; this study
B H ₂ O SE position (J2000.0).....	$3^{\text{h}}29^{\text{m}}12.0104^{\text{s}}, +31^{\circ}13'07.955''$	VLA, 1998 Aug 12; this study
Peak flux (2.7 mm).....	58 mJy	(2'') Looney et al. (2000)
Peak flux (3.6 cm).....	0.33 mJy	Reipurth et al. (2002)
Bolometric luminosity	$14 L_{\odot}$	André et al. (2000)
Bolometric temperature.....	36 K	André et al. (2000)

The synthesized beam was 300 by 275 mas at a position angle of -83° . The strong masers in SSV13, some $196''$ to the north-northwest, contributed sidelobe emission in some spectral channels. In general, however, those channels affected did not have maser emission from IRAS 4A or 4B. With this angular resolution, many spectral features mapped to the same position, although two groups each of spatially separated masers were seen toward both IRAS 4A and 4B. The positions of these four maser groups, as estimated with Gaussian fits to the strongest maser in each group using the AIPS task JMFIT, are listed in Table 1.

For the continuum observations, after the standard calibration, the narrowband data were imaged and used in an iterative self-calibration procedure. The self-calibration solutions obtained were applied to the broadband data, which were imaged with natural weighting of the u, v data (which increases the sensitivity relative to uniform weighting). The synthesized beam for this mode was 360 by 325 mas at a position angle of 77° . The rms noise in this image was $\sim 120 \mu\text{Jy beam}^{-1}$. Unfortunately, the weather was not particularly good, and the system temperatures were elevated by about a factor 1.5–2.0 over good weather conditions. Integrated flux densities at $\lambda = 1.3 \text{ cm}$ for IRAS 4A1 and 4BW were 1.33 and 0.54 mJy, respectively. IRAS 4A2 and 4BE were not detected to a 3σ peak flux density limit of $0.36 \text{ mJy beam}^{-1}$.

2.2. VLBA Observations

We have used the VLBA to obtain four epochs of observation of the water maser emission associated with a number of low-mass YSOs. Two subsequent papers will present results for IRAS 16293–2422 and SVS 13. A previous paper (Claussen et al. 1998) presented results for IRAS 05413–0104. Preliminary results of the observations presented here have appeared in Marvel et al. (2002), while those for IRAS 16293–2422 and SVS 13 have appeared in Wootten et al. (1999, 2002).

The observations were made with one polarization pair (both right- and left-circular polarization) of baseband channels, each with 4 MHz of bandwidth. This provides 54 km s^{-1} of total velocity coverage. A correlation was made with 256 spectral channels per baseband channel, yielding a velocity width in each spectral channel of 0.21 km s^{-1} . Although the velocity coverage is then only $\pm 27 \text{ km s}^{-1}$ from the systemic, no published higher velocity water masers have been detected from the IRAS 4 sources (Claussen et al. 1996; Brand et al. 2003; Furuya et al. 2003).

Since IRAS 4A and 4B are located close together on the sky (within the primary beam of the VLBA antennas), we were able to direct the antennas of the VLBA to a position between the two sources and then perform two separate correlation passes with the VLBA correlator at the precise positions of the sources without loss of sensitivity. The correlation positions coincide with the centimeter continuum positions (see Table 1). Since the data were obtained contemporaneously, simply being correlated in separate passes, the two data sets share identical time ranges and midtimes of observation. The four observations were separated in time by approximately 3 weeks. This separation was chosen based on past single-dish monitoring observations (Claussen et al. 1996; Wilking et al. 1994), which indicated significant flux changes on this timescale. The dates of the observations, midtimes of each epoch, and Julian date of the midtimes are given in Table 2, along with the number of days between each pair of the observations.

The data from each epoch were reduced using AIPS. Strong continuum sources were observed during each epoch to calibrate the bandpass response of the receiver systems of each antenna, as well as to determine any residual delays uncompensated for by the correlator model. After residual delay removal, a strong maser feature displaying simple structure was selected for determination of delay changes during the observations. After the removal of these residual delay rates, the same channel was split

TABLE 2
DATES AND TIMES OF OBSERVATIONS

Epoch	Date (1998)	Mid-time (UT)	Julian Date −2,451,000.0	Δ Time (days)
I.....	Sep 15	11:01:23	71.95928	0.00000
II.....	Oct 4	10:01:30	90.91771	18.95843
III.....	Oct 27	08:01:30	113.83437	41.87509
IV.....	Nov 14	06:31:30	131.77187	59.81259

from the data set for iterative self-calibration. The amplitude and phase self-calibration solutions were then applied to all spectral channels, and the data were edited for poorly calibrated amplitudes.

After all calibrations were applied, the data were mapped using the AIPS task IMAGR. Spectral cubes were formed with pixel cell sizes of $85 \mu\text{as}$. The restoring beam varied in size, but was typically about $800 \mu\text{as}$ by $400 \mu\text{as}$, with a position angle very close to -10.0° east of north. The resolution varied for each epoch depending on the details of the self-calibration process and whether or not good self-calibration solutions were obtained for given antennas at given times. The exact imaging parameters for each epoch and rms noises for typical channels (those channels whose noise floor is not dictated by dynamic range constraints) are listed in Table 3.

To aid in the analysis of the maser motions observed, two-dimensional Gaussian components were fit to each maser spot in each channel where the emission from an individual spot exceeded 5 times the rms noise for that channel. This was carried out using an automated routine in AIPS (SAD) in an iterative process that first fit the strongest Gaussians in a given channel, subtracted the resulting component from the map, and then reiterated the search for Gaussian components above the noise cutoff in the map. This produces a fairly reliable set of Gaussian components without the very tedious process of fitting each component above the limit in each channel manually. The resulting positions were then grouped together in velocity space and position in tabular form. Only masers detected in two or more channels above the cutoff and no more than one beamwidth offset in position were retained. Fewer than 10% of the components were excluded in this process, pointing to the reliability of the iterative fitting process.

The surviving component fits were then visually inspected using the three-dimensional data analysis tool XGOBI.³ Using XGOBI, it is possible to excise spurious Gaussian fits that result from the automated Gaussian fitting routine we utilized. If a fit was called into question by the XGOBI inspection process, the original maps were inspected prior to deleting the data point. Usually, spurious fits were caused by residual sidelobes being fit in several channels by the SAD routine in AIPS and occurred in channels that were dynamic-range limited. The components were then used for all subsequent fitting and analysis as described in the Appendix.

3. DESCRIPTION AND ANALYSIS OF THE MASER REGIONS

3.1. IRAS 4A

At the time of our observations, the water masers of IRAS 4A were found in two dominant regions, separated by 251 mas

TABLE 3
IMAGING PARAMETERS FOR EACH EPOCH

Epoch	Beam Dimensions (μas)	Beam P.A. (deg E of N)	rms (mJy beam $^{-1}$)
4A-I.....	1260×910	11.8	10.4
4A-II.....	830×440	−21.2	6.8
4A-III.....	810×400	−19.1	11.7
4A-IV.....	650×370	−8.2	9.6
4B-I.....	650×350	−14.4	13.9
4B-II.....	620×330	−12.4	7.6
4B-III.....	590×330	−10.8	9.9
4B-IV.....	720×420	−5.8	6.0

(epoch I) or 80.3 AU. A line connecting the two maser regions lies at a position angle of about -50° (north through east), roughly perpendicular to the CO outflow axis. Comparison with our VLA observations show that the masers are associated with IRAS 4A2. Figure 1 shows the spatial and kinematic distribution of the masers for all four epochs, each with a close-up of the southeast (SE) maser component.

The northwest (NW) region exhibited only a single maser spot. Its peak flux decreased in time from 0.97 Jy at the first epoch to 0.12 Jy at the final epoch. The maser is at or slightly redshifted relative to the molecular core velocity of $+6.7 \text{ km s}^{-1}$. The SE region contains more highly redshifted maser spots in epochs I and II, including the maser used to self-calibrate each epoch and another weaker maser just west of the reference spot. In epochs III and IV, another maser appeared ~ 10 mas north of the reference feature at a velocity of $\sim 12.5 \text{ km s}^{-1}$ (Fig. 1). Because we self-calibrated the data using the brightest spot as the reference, its position remains constant over time and appears at the origin.

The separation vector from the reference feature to the NW feature increased in magnitude over time, expanding from 80.4 AU in the first epoch to 82.3 AU in the final epoch. The projected speed of expansion is 53.0 km s^{-1} , given the assumed distance of 320 pc to the source. Since the line-of-sight velocity difference between the two regions is only 2 km s^{-1} , and assuming that this velocity difference and the proper motion result from the projected velocity of an outflow jet, we can calculate an inclination for the outflow of 2° . The true space velocity is therefore nearly equal to that measured in the plane of the sky, namely 53 km s^{-1} . As shown in Table 4, the proper-motion vectors for the NW feature relative to the reference feature has a position angle of -49.8° , identical to the position angle of the separation vector.

3.2. IRAS 4B

Similar to IRAS 4A, IRAS 4B exhibits water maser emission in two dominant well-separated regions associated with IRAS 4BW: one to the NW and another to the SE. Masers in the NW region are redshifted relative to the cloud velocity of 7 km s^{-1} , with velocities ranging from 13 to 20 km s^{-1} . Masers in the SE regions are blueshifted, with velocities ranging from -3 to 4 km s^{-1} . The radial velocities for the NW and SE groups are consistent with those in the northern and southern lobes of the bipolar HCN outflow observed by Choi (2001) and the bipolar H_2CO outflow mapped by Di Francesco et al. (2001). A complete list of the maser component velocities, fluxes, and positions can be found in the Appendix.

Figures 2 and 3 show the overall spatial and kinematic distribution of the maser emission as observed in each epoch for the two maser groups in IRAS 4B. Over time, the masers in the NW

³ XGOBI, a multivariate data analysis tool that operates under the X11 system is freely available for download at <http://www.research.att.com/areas/stat/xgobi/>.

FIG. 1.—Spatial and kinematic distribution of the water masers associated with IRAS 4A2 for the four epochs observed. The top frame in each panel shows the velocity and flux of the detected features, and the lower frame shows the spatial distribution. The masers to the NW are colored blue (at the cloud velocity) and those to the SE are coded red (redshifted relative to the cloud velocity). The size of the circles in the figure is proportional to the flux of that maser component. All components measured in each channel are presented (e.g., no component averaging has been performed). The error in the position of the spots (typically less than $20 \mu\text{as}$) cannot be plotted on the large-scale figure.

expand away from the reference feature at (0, 0) in the SE maser group. A line connecting the two regions lies at a position angle of roughly -29° , and the two regions are separated by 488 mas, or 156 AU. The SE grouping is roughly arclike, with an overall extent of 38 mas (12 AU) and a thickness of 2 mas (0.6 AU), while the NW clump is more linear and extends about 26 mas (8.3 AU), with a thickness roughly the same as the SE clump. The highest velocity maser emission lies at the northern tip of the NW group ($+20 \text{ km s}^{-1}$) and at the southern tip of the SE arc (-3 km s^{-1}). In addition, a small maser group is located approximately 100 mas (32 AU) to the NNW of the main SE arc. This region's structure changed during the epochs, but usually consisted of two or three

well-separated maser features (by several mas, or about 0.3 AU). This distinct region is located close to the line connecting the NW and SW maser groups. If the active maser regions mark the interaction of a bipolar jet from IRAS 4BW with ambient gas, then the opening angle of the jet would be $\sim 10^\circ$.

Beginning in the second epoch, but much more clearly shown in the third and fourth epochs, a new arclike ridge of lower velocity emission became visible to the south of the other masers in the NW region. This region is not well-characterized by pointlike emission, but is more diffuse and seems to be an extended ridge of maser emission. By the fourth epoch, the entire NW feature had grown to roughly 40 mas (12.8 AU) in extent.

TABLE 4
PROPER MOTIONS FOR MASERS OBSERVED IN ALL FOUR EPOCHS

Epoch ID	Velocity (km s ⁻¹)	Peak Flux (Jy)	X Offset (mas)	X _{ref} (mas)	Y Offset (mas)	Y _{ref} (mas)
IRAS 4A Reference						
I-A-REF ^a	10.40	3.29	0.083	...	0.073	...
II-A-REF.....	10.55	4.14	-0.063	...	-0.185	...
III-A-REF.....	10.68	7.97	0.007	...	-0.030	...
IV-A-REF.....	10.72	14.9	-0.005	...	0.002	...
IRAS 4A 1: $v_{\text{exp}} = 53.02$, P.A. = -49.82°						
I-A-1.....	7.61	0.969	-167.543	-167.626	187.223	187.150
II-A-1.....	7.65	0.211	-169.013	-168.950	187.985	188.170
III-A-1.....	7.71	0.208	-170.631	-170.638	189.606	189.636
IV-A-1.....	7.85	0.123	-172.098	-172.093	190.924	190.921
IRAS 4A 2 ^b						
I-A-2.....	10.34	1.90	-1.859	-1.942	-0.086	-0.160
II-A-2.....	10.36	2.26	-1.906	-1.844	-0.118	0.067
III-A-2.....	10.29	1.74	-1.844	-1.851	0.695	0.725
IV-A-2.....	10.35	1.27	-1.938	-1.933	1.113	1.108
IRAS 4B Reference						
I-B-REF ^c	-0.750	36.8	-0.020	122.668	0.0782	-227.845
II-B-REF.....	-1.161	20.4	-0.001	123.598	-0.041	-228.993
III-B-REF.....	-0.224	56.6	-0.000	124.475	-0.013	-229.284
IV-B-REF.....	-0.273	34.3	-0.002	125.498	0.010	-230.661
IRAS 4B 1: $v_{\text{exp}} = 45.40$, P.A. = 144.12°						
I-B-1.....	0.380	0.368	-1.736	120.932	10.016	-217.830
II-B-1.....	0.302	2.14	-1.677	121.921	9.384	-219.549
III-B-1.....	0.485	0.256	-1.626	122.849	8.726	-220.559
IV-B-1.....	-0.158	0.09	-1.732	123.767	8.912	-221.749
IRAS 4B 2: $v_{\text{exp}} = 41.03$, P.A. = 140.07°						
I-B-2.....	2.81	0.829	-9.429	113.239	33.556	-194.290
II-B-2.....	3.05	0.134	-9.531	114.067	33.844	-195.089
III-B-2.....	3.25	1.24	-9.074	115.402	32.605	-196.679
IV-B-2.....	3.39	0.218	-9.417	116.082	32.975	-197.686
IRAS 4B 3: $v_{\text{exp}} = 41.90$, P.A. = 40.5°						
I-B-3.....	17.46	0.634	-239.832	-117.165	432.737	204.892
II-B-3.....	17.07	2.80	-241.622	-118.024	434.953	206.020
III-B-3.....	17.27	0.521	-243.675	-119.200	436.448	207.164
IV-B-3.....	17.13	0.318	-245.662	-120.163	439.063	208.402
IRAS 4B 4: $v_{\text{exp}} = 41.47$, P.A. = -36.5°						
I-B-4.....	17.74	2.31	-239.610	-116.943	430.618	202.772
II-B-4.....	17.66	5.90	-241.452	-117.853	432.924	203.991
III-B-4.....	17.60	1.50	-243.359	-118.884	434.451	205.167
IV-B-4.....	17.43	0.862	-245.111	-119.612	437.038	206.377
IRAS 4B 5: $v_{\text{exp}} = 46.15$, P.A. = -36.44°						
I-B-5.....	17.95	0.643	-239.816	-117.148	438.361	210.516
II-B-5.....	17.99	0.957	-241.704	-118.106	440.878	211.947
III-B-5.....	17.84	1.24	-243.767	-119.292	442.811	213.527
IV-B-5.....	17.22	1.55	-245.495	-119.997	445.035	214.374

^a Strongest maser in IRAS 4A used as both a position and velocity reference.

^b Component A-2 is so close to A-REF that v_{exp} is not meaningful in this case.

^c Strongest maser in IRAS 4B used as a position reference only. Proper motions are measured relative to a fiducial point for each epoch given as X_{ref} and Y_{ref} for this maser (see text).

FIG. 2.—Spatial and kinematic distribution of the blueshifted masers southeast of IRAS 4BW for all four epochs. In each panel, the top frame shows the velocity (color-coded) and flux (symbol size) of the detected features, and the lower frame shows the spatial distribution and the velocity distribution as a function of declination. The uncertainty in the fitted position of each feature cannot be displayed at this scale, but in the worst case is of the order of $20 \mu\text{as}$.

We compared the positions of the reference feature, which is present in all epochs in the SE clump (B-REF in Table 4) with a maser in the NW clump that persisted through all four epochs (B-5 in Table 4). The separation vector for the two masers increased over 59.81262 days from 499.59 mas (159.87 AU) to 508.25 mas (162.64 AU). A linear regression fit to the magnitude of the separation vector as a function of time yields an expansion velocity of 78.2 km s^{-1} , with a formal error of $\pm 3.9 \text{ km s}^{-1}$. The radial velocity difference between the two masers of 18.5 km s^{-1} suggests an inclination of the outflow of about 13° from the plane of the sky.

To investigate the proper motions of the masers associated with IRAS 4BW, we identified all components that were present

in all four epochs of observation. Three components were identified in the NW, and two were identified in the SE. We averaged the mean positions of the NW and SE components and averaged these two positions to estimate an average geometric fiducial position for each epoch. This position was then subtracted from each maser component. These referenced positions are presented along with the fluxes and original positions in Table 4. Finally, we fit straight lines to the positions of the components as a function of time. The resulting tangential velocities and position angles from these fits are also presented in Table 4 and are very similar for all five maser components. Their motions are very uniform over the four epochs, with average values of $43 \pm 2 \text{ km s}^{-1}$ at a position angle of $-38^\circ \pm 2^\circ$. Figure 4 shows all of the maser

FIG. 3.—Spatial and kinematic distribution of the redshifted masers northwest of IRAS 4BW for all four epochs. In each panel, the top frame shows the velocity (color-coded) and flux (symbol size) of the detected features, and the lower frame shows the spatial distribution and the velocity distribution as a function of declination. The uncertainty in the position of each feature cannot be displayed at this scale, but in the worst case is of the order of $20 \mu\text{as}$.

components as observed in epoch 1 and their corresponding proper-motion vectors. Again, an overall expansion is apparent along the axis connecting the two regions.

4. DISCUSSION

4.1. IRAS 4A

Since the water masers toward IRAS 4A have always been found coincident with the position of A2 (when they were detected), we consider three alternatives for the origin of the water maser emission seen using the VLBA: (1) the water masers are in an outflow that is the close-in extension of the large-scale molec-

ular outflow, driven by IRAS 4A2; (2) the water masers are excited by an interaction of the close-in flow and the accretion disk itself; and (3) the water maser outflow is driven by a companion to IRAS 4A2, as yet unresolved. We consider each of these alternatives in turn.

4.1.1. Alignment with the Large-Scale Outflow

The alignment of the water maser outflow is curious in that it does not have the same position angle as any of the larger-scale outflows emanating from the two sources in IRAS 4A. The axis of the large-scale outflow defined by the SiO, CS, CO, and H₂ emission (Choi et al. 2005, 2006; Blake et al. 1995), and driven by A2,

FIG. 4.—Plot of the proper motions for the five maser components present in all four epochs for IRAS 4BW relative to their positions in epoch 1 (see Table 4). The length of the proper motion vectors are 2 km s^{-1} per plot unit. Note that the proper motions are not at the same position angle as a line connecting the two maser groups (see text).

may be drifting with a rate of $0.011^\circ \text{ yr}^{-1}$ (counterclockwise as seen on the sky) and a zero-point angle of around -166° (i.e., presumably the jet-injection angle), depending on the fit of the large-scale emission (see Fig. 3a of Choi et al. 2006). For the large-scale flow, the blueshifted emission is generally in the south to southwest, and the redshifted emission is generally in the north and northeast. The maser outflow, however, on a scale of 200 mas, has a position angle of -50° , and has the more highly redshifted emission to the southeast. Thus, if the masers were marking the close-in extension to the flow, the position angle of the flow would have to change from $+130^\circ$ to -165° , in a counterclockwise direction, from 40 to 300 AU from the star. We consider that such a shift in jet direction is very unlikely, if not unphysical.

4.1.2. Interaction with the A2 Accretion Disk

The jet direction near the base of the flow, as indicated by Choi et al. (2006) and also by the $\text{CS } J = 7 \rightarrow 6$ emission (Blake et al. 1995) appears to be about 0° (i.e., north-south). If the masers were excited by an interaction of the outflow with the A2 accretion disk itself, then the opening half-angle of the jet at its base would have to encompass the direction of the maser axis. Thus, the opening half-angle of the jet would be $\sim 50^\circ$. Such a large opening angle would seem inconsistent with the Class 0 phase and the high degree of collimation seen in the larger scale outflow.

4.1.3. An Unresolved Companion to IRAS 4A2

Perhaps the simplest explanation for the direction of the axis of the water maser emission is that it is not related directly at all

to the large-scale molecular outflow, but that the masers are located in an independent outflow from another young stellar object. A large fraction of the water masers around low-luminosity YSOs, which have been observed and mapped with very high angular resolution, show bipolar flows and structural evidence for bow shocks as the jet impinges on and sweeps up ambient molecular material. Although the structural evidence is sparse (because of the small number of maser spots) for IRAS 4A, the bipolar velocity separation strongly suggests that the masers are tracing a jet outflow as, for example, in IRAS 05413–0104 (Claussen et al. 1998). It is clear from VLA observations (e.g., Furuya et al. 2003; this paper) that the water masers toward IRAS 4A are always found toward A2 rather than A1, so A1 cannot drive the maser outflow. Therefore, we suggest that in this scenario, there must be a companion to A2 that drives the water masers. Since the masers are clearly associated with the position of A2, and not A1, it is clear that the maser outflow, in this scenario, is *not* associated with A1. Indeed, the presence of a companion to A2 separated by 30–80 AU has been proposed by Choi et al. (2006) to explain the drifting of the 4A2 outflow axis. A companion within 100 mas (32 AU) of A2 would be consistent with it not being resolved in $\lambda = 3.6 \text{ cm}$ continuum observations with the VLA (Reipurth 2000).

4.2. IRAS 4B

While the masers trace a jet similar in velocity to the HCN outflow, the position angle of -29° is significantly different from the north-south orientation of the outflow. Choi (2001) suggests a dynamical lifetime for the molecular outflow of $\sim 200 \text{ yr}$ (corrected to a distance of 320 pc). This is much longer than the dynamical lifetime of the maser emission, which is on the order of a year. We conclude that the jet must be precessing in such a way as to move the projected outflow axis from a position angle of 0° to -29° in several hundred years. In this picture, the jet would be impacting the walls of the outflow cavity as it precessed “clockwise” in the plane of the sky. The NW maser group would represent the western wall of the preexisting cavity, and the SE maser group the eastern wall. This would explain the linear, arclike structure of the maser groups and the appearance of new masers at the southern extent of the NW group. This model can only be reconciled with observations of dense outflowing gas traced in H_2CO by Di Francesco et al. (2001) if there were a second outflow from an unresolved companion to IRAS 4BW. If our picture is correct, one would expect future VLBI observations to reveal new masers south (north) of the NW (SE) group. One would also expect that on a larger scale, tracers of the jet would display a characteristic S-shape, bending back to a north-south orientation. Indeed, slowly precessing jets have been proposed for a number of YSO outflows (Matthews et al. 2006; Hodapp et al. 2005; Terquem et al. 1999) that suggest the presence of a nearby companion (Terquem 1998).

The proper-motion vectors for the five persistent maser spots relative to a fiducial point (Fig. 4) suggest a larger position angle than that of the outflow axis defined by the relative motion of SE and NW masers (-38° vs. -29° ; see Table 4). This is reminiscent of the proper motions seen in the triple radio continuum source in Serpens (Curiel et al. 1993), where the outer components of the radio jet have the same tangential velocity but are moving in slightly different directions. Although the maser emission in IRAS 4B traces the jet at scales 40 times smaller than that traced by the radio continuum in the Serpens jet, this difference in relative motions may be suggestive of similar physical processes in the two jets.

The larger position angle defined by the proper motion vectors in IRAS 4B suggest that, in addition to the tangential velocity imparted to the masing gas by the outflow jet at a position angle of -29° , there is another component of about 6 km s^{-1} perpendicular to the outflow axis. It is possible that this motion is related to the expansion of the outflow cavity into the ambient cloud.

4.3. Maser Geometry

We estimate that the maser outflows in both IRAS 4A and 4B are nearly in the plane of the sky (inclination of only 2° for IRAS 4A and about 13° for IRAS 4B). In Claussen et al. (1998) we estimated that the inclination of the maser outflow for IRAS 05413–0104 was only 4° . Models of maser emission (e.g., Hollenbach 1997) suggest that the masers trace shocks produced by the interaction of the jet and ambient material along the jet. Shocks that propagate close to the plane of the sky provide much longer maser coherence lengths than those that propagate at larger inclination angles. Thus, outflows that have inclinations close to the plane of the sky provide a more favorable maser geometry, perhaps explaining why some well-known outflows from low-mass YSOs do not show maser emission.

5. CONCLUSIONS

We have observed the water masers associated with IRAS 4A and 4B at VLBI resolutions in four epochs over three months. We have determined that the masers are related to the jets emanating from these YSOs due to their spatio-kinematic distribu-

tion. In both sources the masers are found to be associated with known single components of multiple systems, with total separation velocities between 53 and 78 km s^{-1} .

This is further confirmed by the large proper motions measured for both sources, which clearly rules out rotation due to the mass constraints placed on the central objects by other observations. The water masers of IRAS 4B form arclike structures roughly 10 AU in length and less than 0.6 AU in thickness. The structure of these structures changes rapidly with time, with many new maser components appearing and disappearing in just 1 month. The orientation of the structures in the plane of the sky does not agree with larger scale outflow angle, which we attribute to a possible unseen very close companion. Future observations will have to sample the source more frequently than once every 3 weeks, with once every 3–5 days probably providing the best results.

K. B. M. wishes to thank the AAS Council and AAS Executive Officer, Robert Milkey, for allowing him to continue research work in his current position. K. B. M. thanks T. E. K. for her continued support. And special thanks to T. Beasley for exposing K. B. M. to the rigors of Outback observing and a grueling exercise regimen, during which time a portion of this paper was written. Finally, K. B. M. thanks the NRAO for hosting him during a research leave from the AAS, during which time the bulk of this paper was completed.

APPENDIX

TABLE OF MASER COMPONENTS FOR IRAS 4B

Maser component fits were averaged using a flux-squared weighting scheme, which emphasizes the strongest portion of the maser emission and produces a reliable location for the maser. As some masers are present for only three channels, a full Gaussian fit to the maser as a function of velocity was deemed impractical, although it could have been used for the masers that existed over many channels. We opted to use the same position-determining scheme for all features. These averaged component fits and their corresponding errors are given in Table 5 for IRAS 4B. Note that the uncertainty in position for the reference channel is simply the error on the mean position values determined for the spot and is indicative of the reliability of the Gaussian fitting routines in the presence of noise in the images.

TABLE 5
FITTED GAUSSIAN COMPONENT IN IRAS 4B

Epoch ID	Other	Velocity (km s^{-1})	Peak Flux (Jy)	X Offset (mas)	Y Offset (mas)	σ_X (μas)	σ_Y (μas)
I-SE-1	B-2	2.81	0.83	−9.428	33.555	4.4	8.0
I-SE-2		2.30	8.24	−7.362	27.891	0.1	1.1
I-SE-3		2.21	4.97	−5.722	21.858	1.1	1.3
I-SE-4		2.27	0.76	−5.966	22.825	5.5	5.3
I-SE-5		2.10	0.44	−6.433	23.448	7.8	12.9
I-SE-6		1.94	0.76	−5.327	20.908	3.8	5.5
I-SE-7		1.88	1.76	−5.883	22.016	1.2	2.5
I-SE-8		1.74	0.58	−5.247	20.739	6.1	11.8
I-SE-9		1.33	0.37	−4.515	18.473	7.6	14.6
I-SE-10		1.29	0.38	−3.104	15.207	8.1	14.7
I-SE-11	B-1	0.38	0.37	−1.735	10.010	6.6	8.8
I-SE-12		0.23	0.62	−1.442	8.355	3.7	2.9
I-SE-13		0.16	5.84	−1.571	9.156	0.3	1.2
I-SE-14		0.10	0.41	−1.281	6.440	6.7	12.1
I-SE-15	B-REF	−0.01	0.48	−1.392	7.743	7.2	21.2
I-SE-16		−0.75	36.76	−0.020360	0.07815	0.007	0.04
I-SE-17		−1.57	0.20	−49.551	88.150	14.0	25.6
I-SE-18		−2.68	0.38	−48.231	84.987	7.42	13.1
I-NW-1	B-5	17.95	0.64	−239.816	438.361	4.8	9.4
I-NW-2	B-4	17.74	2.31	−239.610	430.618	1.4	3.5
I-NW-3		17.71	0.39	−239.404	429.748	6.2	4.5

TABLE 5—*Continued*

Epoch ID	Other	Velocity (km s ⁻¹)	Peak Flux (Jy)	<i>X</i> Offset (mas)	<i>Y</i> Offset (mas)	σ_X (μ as)	σ_Y (μ as)
I-NW-4	B-3	17.70	0.27	−239.354	427.401	8.8	23.2
I-NW-5		17.46	0.63	−239.832	432.737	5.1	2.2
I-NW-6		17.34	8.64	−239.793	431.960	0.1	1.0
I-NW-7		17.31	0.52	−239.674	431.152	4.5	2.4
I-NW-8		17.13	2.13	−239.178	420.144	1.6	5.1
I-NW-9		16.93	3.15	−239.996	425.710	1.3	3.3
I-NW-10		16.88	0.72	−240.000	424.172	5.2	22.2
I-NW-11		16.71	0.46	−239.842	421.856	7.3	24.4
I-NW-12		16.28	0.23	−239.074	413.088	14.8	31.3
I-NW-13		16.11	0.12	−239.817	415.178	28.7	54.5
II-SE-1	B-2	3.08	0.13	−9.929	34.227	18	28
II-SE-2		3.05	0.13	−9.532	33.844	18	23
II-SE-3		2.53	0.32	−6.009	22.195	6	10
II-SE-4		2.38	1.09	−7.803	28.253	2	3
II-SE-5		2.26	0.46	−6.057	22.259	5	7
II-SE-6		2.24	0.23	−6.545	23.227	10	17
II-SE-7		1.63	0.18	−2.539	15.374	11	21
II-SE-8		1.60	0.48	−4.945	19.856	5	9
I-SE-9		1.50	0.64	−4.899	19.603	3	1
II-SE-10		1.48	1.23	−4.751	18.793	2	5
II-SE-11	B-1	1.22	0.25	−3.393	15.304	10	17
II-SE-12		1.13	0.14	−2.706	13.848	19	30
II-SE-13		1.10	0.17	−2.724	13.181	10	22
II-SE-14		0.84	0.53	−2.317	11.923	4	10
II-SE-15		0.43	2.44	−1.758	9.760	1	1
II-SE-16		0.43	1.05	−1.558	8.746	1	1
II-SE-17		0.30	2.14	−1.677	9.383	1	2
II-SE-18		0.21	1.03	−1.298	6.382	2	4
II-SE-19		0.11	0.29	−0.580	2.781	6	13
II-SE-20		0.05	0.70	−1.401	7.519	2	5
II-SE-21	B-REF	−0.97	12.80	−0.132670	0.679741	0.002	0.005
II-SE-22		−1.16	20.41	−0.000980	−0.041125	0.004	0.007
II-SE-23		−1.52	2.11	0.244	−1.140	1	1
II-SE-24		−1.90	0.32	−46.0974	77.935	6	11
II-SE-25		−2.04	0.24	1.360	−6.688	6	11
II-SE-26		−2.37	0.29	−47.5718	82.438	7	17
II-NW-1	B-5	17.99	0.96	−241.704	440.878	2	5
II-NW-2		17.99	0.86	−241.869	449.662	2	4
II-NW-3	B-4	17.66	5.90	−241.452	432.924	1	1
II-NW-4		17.61	1.48	−241.606	433.736	2	10
II-NW-5	B-3	17.10	0.34	−241.901	429.900	9	38
II-NW-6		17.09	0.56	−241.571	427.050	5	5
II-NW-7		17.07	2.80	−241.622	434.953	1	1
II-NW-8		17.04	3.58	−241.631	428.024	1	2
II-NW-9		17.03	0.26	−241.059	423.332	11	6
II-NW-10		17.02	3.29	−241.073	422.442	1	1
II-NW-11		16.97	0.37	−241.235	421.489	6	7
II-NW-12		16.95	0.26	−241.421	424.037	9	28
II-NW-13		16.90	0.43	−241.797	429.376	6	12
II-NW-14		16.87	0.54	−241.683	426.113	5	19
III-SE-1	B-2	3.25	1.24	−9.074	32.605	2	2
III-SE-2		3.19	0.32	−4.623	22.469	6	11
III-SE-3		1.83	0.18	−2.942	16.312	11	17
III-SE-4		1.65	0.18	−4.944	18.837	9	23
III-SE-5		1.38	0.14	−4.707	17.604	12	22
III-SE-6		0.99	0.81	−2.259	10.991	3	5
III-SE-7		0.60	0.18	−1.151	6.024	10	22
III-SE-8		0.51	0.12	−1.388	7.749	13	23
III-SE-9	B-1	0.48	0.26	−1.626	8.726	8	18
III-SE-10		0.30	0.16	−1.118	5.494	9	28
III-SE-11	B-REF	0.09	0.73	−0.421	1.823	3	7
III-SE-12		−0.22	56.60	−0.000320	−0.013123	0.001	0.005
III-SE-13		−0.26	2.17	0.134	−0.762	1	1

TABLE 5—*Continued*

Epoch ID	Other	Velocity (km s ⁻¹)	Peak Flux (Jy)	X Offset (mas)	Y Offset (mas)	σ_X (μ as)	σ_Y (μ as)
III-SE-14		-0.67	0.22	1.200	-4.693	11	21
III-SE-15		-0.84	2.92	-45.799	75.965	1	1
III-SE-16		-0.97	0.16	2.054	-8.797	9	29
III-SE-17		-1.01	0.13	0.504	-2.326	14	21
III-SE-18		-1.34	1.33	1.165	-5.231	1	3
III-SE-19		-1.34	0.23	0.947	-4.513	7	9
III-SE-20		-1.69	0.20	-45.966	76.791	10	11
III-SE-21		-2.72	0.39	-48.936	84.458	6	10
III-NW-1		17.99	3.58	-243.777	451.560	0.3	1
III-NW-2	B-5	17.84	1.24	-243.767	442.811	2	4
III-NW-3		17.64	0.52	-243.826	442.337	5	10
III-NW-4	B-4	17.60	1.50	-243.359	434.451	1	3
III-NW-5		17.34	0.99	-244.031	441.137	2	4
III-NW-6		17.28	1.26	-243.607	435.704	2	4
III-NW-7	B-3	17.27	0.52	-243.675	436.448	4	6
III-NW-8		17.14	0.31	-243.397	430.521	8	23
III-NW-9		17.02	1.00	-243.439	429.509	3	5
III-NW-10		16.94	0.60	-243.038	424.096	3	6
III-NW-11		16.86	0.14	-243.471	427.598	13	33
III-NW-12		16.83	0.30	-243.662	431.119	10	24
III-NW-13		16.49	0.39	-244.642	437.123	6	13
III-NW-14		16.20	0.10	-244.699	434.707	25	84
IV-SE-1	B-2	3.39	0.22	-9.417	32.975	7	12
IV-SE-2		2.30	0.48	-0.527	10.169	4	7
IV-SE-3		1.18	0.08	1.014	2.647	22	36
IV-SE-4	B-1	-0.16	0.09	-1.732	8.912	17	25
IV-SE-5	B-REF	-0.27	34.28	-0.001790	0.009699	0.001	0.004
IV-SE-6		-1.46	0.09	2.713	-10.796	19	33
IV-SE-7		-1.52	0.31	-46.019	76.637	5	10
IV-SE-8		-1.65	0.07	-46.833	78.663	24	5
IV-NW-1		19.07	0.19	-245.612	458.674	7	13
IV-NW-2		18.99	2.04	-245.643	456.618	1	1
IV-NW-3		18.80	0.48	-245.576	458.611	3	14
IV-NW-4		18.60	0.47	-245.375	460.263	3	5
IV-NW-5		18.29	0.30	-245.403	453.824	4	8
IV-NW-6		17.81	0.59	-245.703	449.530	3	6
IV-NW-7	B-4	17.43	0.86	-245.111	437.038	2	6
IV-NW-8	B-5	17.22	1.55	-245.495	445.035	1	2
IV-NW-9		17.13	0.92	-245.701	443.219	2	4
IV-NW-10	B-3	17.13	0.32	-245.662	439.063	9	14
IV-NW-11		17.09	0.35	-245.643	444.135	6	8
IV-NW-12		17.05	0.61	-246.666	446.387	3	8
IV-NW-13		16.27	0.14	-246.060	440.124	15	16
IV-NW-14		16.14	0.99	-246.162	439.176	2	4
IV-NW-15		14.33	0.52	-247.708	425.411	2	7

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