

**San Jose State University**

---

**From the Selected Works of Aaron J. Romanowsky**

---

September 11, 2017

# The SLUGGS Survey: trails of SLUGGS galaxies in a modified spin- ellipticity diagram

Sabine Bellstedt, *Swinburne University of Technology*

Alister W. Graham, *Swinburne University of Technology*

Duncan A. Forbes, *Swinburne University of Technology*

Aaron J. Romanowsky, *San Jose State University*

Jean P. Brodie, *University of California Observatories*, et al.



Available at: <https://works.bepress.com/>

aaron\_romanowsky/132/

# The SLUGGS Survey: trails of SLUGGS galaxies in a modified spin-ellipticity diagram

Sabine Bellstedt,<sup>1</sup>★ Alister W. Graham,<sup>1</sup> Duncan A. Forbes,<sup>1</sup>  
Aaron J. Romanowsky,<sup>2,3</sup> Jean P. Brodie<sup>3</sup> and Jay Strader<sup>4</sup>

<sup>1</sup>Centre for Astrophysics and Supercomputing, Swinburne University of Technology, Hawthorn VIC 3122, Australia

<sup>2</sup>Department of Physics and Astronomy, San José State University, One Washington Square, San Jose, CA 95192, USA

<sup>3</sup>University of California Observatories, 1156 High Street, Santa Cruz, CA 95064, USA

<sup>4</sup>Department of Physics and Astronomy, Michigan State University, East Lansing, Michigan 48824, USA

Accepted 2017 May 26. Received 2017 May 21; in original form 2017 April 2

## ABSTRACT

We present radial tracks for four early-type galaxies with embedded intermediate-scale discs in a modified spin-ellipticity diagram. Here, each galaxy’s spin and ellipticity profiles are shown as a radial track, as opposed to a single, flux-weighted aperture-dependent value as is common in the literature. The use of a single ellipticity and spin parameter is inadequate to capture the basic nature of these galaxies, which transition from fast to slow rotation as one moves to larger radii where the disc ceases to dominate. After peaking, the four galaxy’s radial tracks feature a downturn in both ellipticity and spin with increasing radius, differentiating them from elliptical galaxies, and from lenticular galaxies whose discs dominate at large radii. These galaxies are examples of so-called discy elliptical galaxies, which are a morphological hybrid between elliptical (E) and lenticular (S0) galaxies and have been designated ES galaxies. The use of spin-ellipticity tracks provides extra structural information about individual galaxies over a single aperture measure. Such tracks provide a key diagnostic for classifying early-type galaxies, particularly in the era of 2D kinematic (and photometric) data beyond one effective radius.

**Key words:** galaxies: elliptical and lenticular, cD – galaxies: individual: NGC 821, NGC 3377, NGC 4278, NGC 4473 – galaxies: kinematics and dynamics.

## 1 INTRODUCTION

The extent to which rotation is present in galaxies has been a widely explored topic for decades. Rotation profiles for spiral galaxies have been studied since the late 1930s (Babcock 1939), but due to the lower surface brightnesses of the outer regions of early-type galaxies (ETGs), studies of their rotation profiles did not commence until later. A sample of ETGs were observed to be elongated (Sandage, Freeman & Stokes 1970), which was expected to be due to rotation (Larson 1975). Early measurements of rotation profiles in lenticular (S0) galaxies confirmed the expectation of strong rotation (Williams 1975). Astronomers were then surprised by data suggesting that there was little rotation in some elliptical (E) galaxies (e.g. Bertola & Capaccioli 1975, who measured significantly less rotation in the flattened elliptical NGC 4697 than present in spiral galaxies). This unexpected finding led to much research regarding the nature of rotation versus anisotropic velocity dispersion (e.g. Illingworth 1977; Binney 1978; Schechter & Gunn 1979; Davies 1981; Kormendy & Illingworth 1982; Davies et al. 1983).

One of the early manners in which information about rotational support was condensed was the so-called Binney (1978) diagram. This diagram presented the ratio of the maximum rotational velocity to the central velocity dispersion  $v/\sigma$  against the ellipticity  $\epsilon$  for individual galaxies. In particular, this diagram was used by Binney (1978) to address the question of how galaxies that are not rotationally supported can be elongated by anisotropy. Davies et al. (1983) later used this diagram and noted that fainter ellipticals ( $M_B > -20.5$  mag) tended to be rotationally supported, whereas brighter ellipticals ( $M_B < -20.5$  mag) were pressure supported by near-isotropic velocity dispersion.

The use of this diagram was ideal for longslit data, in which a maximum rotation value along the major axis and a central velocity dispersion could be measured. With the increasing use of two-dimensional spectroscopy, in particular via integral field units (IFUs), much more information was available with which to characterize the kinematics of galaxies. In order to utilize 2D measurements of both velocity and velocity dispersion within a specified aperture, an observationally measurable, luminosity-weighted spin parameter, denoted by  $\lambda_R$ , was developed by the SAURON team (Emsellem et al. 2007). The use of spatial weighting allowed this parameter to better indicate the rotation of galaxies with non-conventional kinematics (whose properties could not simply be

★E-mail: sbellstedt@swin.edu.au

summarized by major-axis measurements, such as galaxies with kinematic twists and misalignments as was later discussed by Krajnović et al. (2011). The main outcome of the use of the  $\lambda_R$  parameter by the SAURON (de Zeeuw et al. 2002) and ATLAS<sup>3D</sup> (Cappellari et al. 2011a) teams was to plot this parameter in conjunction with the ellipticity  $\epsilon$  in a spin-ellipticity diagram to identify two classes of early-type galaxies according to their kinematics: fast and slow central rotators.

Despite the use of high-quality 2D kinematic data, the final results are still typically reduced to a single measurement for each galaxy. The breadth of information which can be gained from this diagram is further reduced by the fact that the position in this  $\lambda_R - \epsilon$  space is dependent on the single aperture size, and does not differentiate between galaxies with constant or varying rotation in differing regions within the same galaxy, see figs 2 and 5 in Emsellem et al. (2007). Additionally, ellipticity often changes with radius (as has been noted by numerous studies, e.g. Liller 1966; di Tullio 1978, and references therein).

The use of the  $\lambda_R - \epsilon$  diagram and other kinematic diagnostics has challenged the traditional morphological classification of ETGs into elliptical (E) and lenticular (S0) types, as galaxies previously classified into these morphologies do not form two distinct kinematic populations. One reason for this is that inclination affects disc identification differently for photometric and kinematic techniques.<sup>1</sup> Moreover, photometric decompositions into bulge and disc components often produce inconsistent results requiring the inclusion of 2D kinematic data to help break degeneracies in disc fractions and sizes. The presence of discs in some ETGs was therefore missed in earlier photometric studies. For example, Cappellari et al. (2011b) estimated that only  $\sim 34$  per cent of elliptical galaxies are correctly classified and that the rest contain discs.

The observation of a ‘morphological hybrid’ – an ETG that displays properties somewhere in between those of ellipticals and lenticulars – has been noted in the past (Liller 1966). The presence of these discs confined within the main spheroidal component of the galaxies has been analysed in photometric studies such as that by Scorza & Bender (1995), where the comment was made that there was likely a *continuity of disc properties at the low disc-to-bulge ratio end of the Hubble sequence*. Such embedded discs have also been identified by Cinzano & van der Marel (1994) and Rix & White (1990). Recently, Savorgnan & Graham (2016) and Graham, Ciambur & Savorgnan (2016) discussed the presence of ‘intermediate-scale discs’ in ETGs (i.e. discs that are intermediate in size between nuclear discs and large-scale discs), and their downturning ellipticity and spin profiles. Given that such discs come in a range of sizes, (Nieto, Capaccioli & Held 1988; Simien & Michard 1990; Michard & Marchal 1993; Andreon et al. 1996; Krajnović et al. 2013) the motivation of separating galaxies purely into the binary classification of fast and slow central rotators could be questioned when considering that galaxies may all have fast and slow rotating components of varying proportions (as is evident in the works of, e.g. Arnold et al. 2014; Foster et al. 2016).

Graham et al. (2017) introduced the concept of plotting galaxies as tracks, rather than as single points, in a modified spin-ellipticity diagram, showing their movement as a function of galaxy radius. In this paper, we characterize such radial tracks in  $\lambda(R) - \epsilon(R)$  diagrams for galaxies with intermediate-scale discs and compare

these with tracks of typical E and S0 galaxies. In Section 2, we describe our sample and the data used and in Section 3 we outline our method. The results are given in Section 4, which we discuss in Section 5. We summarize and conclude in Section 6.

## 2 SAMPLE AND DATA

We focus on four galaxies from the SAGES Legacy Unifying Globulars and GalaxieS (SLUGGS) survey<sup>2</sup> (Brodie et al. 2014) noted to have declining spin profiles at large radii by Bellstedt et al. (2017, hereafter B17). These galaxies are NGC 821, NGC 3377, NGC 4278 and NGC 4473. All are classified as elliptical galaxies, with circularised effective radii ( $R_e$ ) of 43.2, 45.4, 28.3 and 30.2 arcsec (4.9, 2.4, 2.1 and 2.2 kpc), respectively. The kinematics for these galaxies behave differently in their outskirts compared to their inner regions, flagging them as potential examples of galaxies with intermediate-scale discs. NGC 4473 had been described by Krajnović et al. (2011) as a ‘double-sigma’ galaxy, characterized by two velocity dispersion peaks along the major axis, exhibiting two distinct inner and outer rotation components (Foster et al. 2013; Alabi et al. 2015). Hence, NGC 4473 does not host an intermediate-scale disc, however we include it to depict the track diagnostic behaviour for such a galaxy. Kinematic maps of all four galaxies using SLUGGS data have been published by Arnold et al. (2014) and Foster et al. (2016). Each of these galaxies displays fast rotation in the central region, which declines at larger radii.

Furthermore, for reference, we present an equivalent analysis for a typical E galaxy (NGC 4365) and a typical S0 galaxy (NGC 1023), and plot radial tracks for other ETGs in the SLUGGS sample. The data used within this study all come from the SLUGGS survey, and were taken with the DEep Imaging Multi-Object Spectrograph on the Keck II telescope in Hawaii. The data reduction procedures are given in previous papers (e.g. Arnold et al. 2014, B17).

## 3 METHOD

Rather than simply plotting an aperture  $\lambda_R$  value against an average ellipticity, as commonly done in the literature, we plot ‘annular’ values of both stellar spin and ellipticity measured at varying radii to produce tracks for individual galaxies. While comparisons of aperture spin measurements of different sizes have been made in the past (see for example, fig. 6 in Raskutti, Greene & Murphy 2014), our approach provides a clearer understanding of the overall kinematic structure of individual ETGs. To do this, we require measurements of both the ellipticity  $\epsilon$  and spin  $\lambda$  profiles. Ellipticity profiles were taken from *Spitzer* data presented in Forbes et al. (2017), and we made new measurements of  $\lambda(R)$  at the corresponding radii from the SLUGGS kinematic data. These ellipticity profiles have been presented in the Appendix for each of the galaxies for which spin–ellipticity tracks have been plotted.

A slightly modified version of the technique outlined in B17 is used to measure  $\lambda(R)$ . At each radius, an annulus of width 2 arcsec is defined with an ellipticity equal to the local isophotal ellipticity of the galaxy, allowing us to measure the local kinematic properties within this annulus. The ellipticity of each annulus varies according to the local measurement, as opposed to a uniform ellipticity value used by B17.

As in B17, 2D maps are produced using the Kriging technique (Pastorello et al. 2014). Kriging produces interpolated 2D maps with an effective smoothing scale of 15–25 arcsec, depending on

<sup>1</sup> From photometry, discs are identified via their projected ellipticity, which scales with inclination  $i$  (where  $i = 0$  for a face-on galaxy) as  $\cos(i)$ , whereas kinematic studies identify discs through rotation, which scales as  $\sin(i)$ .

<sup>2</sup> <http://sluggs.swin.edu.au>

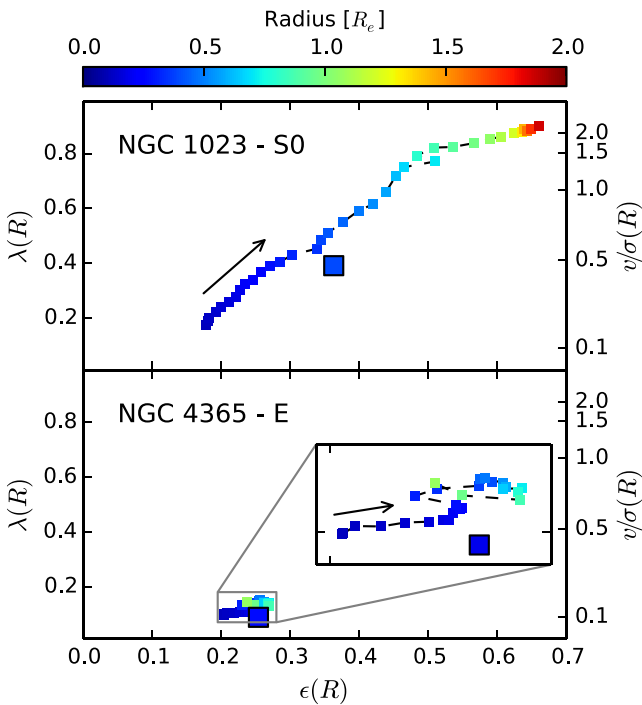
the density of data points. It is the Kriged map, rather than the sparse data points themselves, that we use to calculate kinematic properties. All points in a specified annulus are used to measure both  $\lambda(R)$  and  $v/\sigma(R)$  with the expressions:

$$\lambda(R) \equiv \frac{\sum_k R_k |V_k|}{\sum_k R_k \sqrt{V_k^2 + \sigma_k^2}} \quad \text{and} \quad v/\sigma(R) \equiv \sqrt{\frac{\sum_k V_k^2}{\sum_k \sigma_k^2}}.$$

Here,  $R_k$ ,  $V_k$  and  $\sigma_k$  represent the circularised radius, velocity and velocity dispersion of the  $k$ th Kriging point. The differential flux across a single annulus is neglected. As described by B17, the  $\lambda(R)$  measurement in each annulus is taken as the mean value of 100 bootstrap resamplings of the pixels within each annulus. This process ensures that the variation in pixel number per bin (a result of finite pixel resolution) does not affect the measurement. The spread in  $\lambda(R)$  values measured in each bin is typically 0.001–0.005, indicating that slight kinematic variations across a single annulus do not affect our measurement of the local stellar spin.

We find that  $\lambda(R)$  and  $v/\sigma(R)$  scale as  $\lambda(R) = \kappa v/\sigma(R) / \sqrt{1 + \kappa^2 v/\sigma(R)^2}$  as described by Emsellem et al. (2011), where  $\kappa \simeq 0.9$ , as opposed to  $\kappa \simeq 1.1$  for global values. The minimum radius at which our tracks begin is  $0.1 R_e$ , and the maximum radial extent of our data is defined as the radius where the azimuthal data coverage within the annulus drops to 85 per cent.

To determine whether our relatively sparse spatial sampling of the galaxy kinematics, due to our use of multislit data rather than a contiguous IFU, affects our results, we compare our tracks with those produced by ATLAS<sup>3D</sup> data over the central ( $< 1 R_e$ ) region. We confirm that the results are qualitatively the same.



**Figure 1.** Radial tracks of spin  $\lambda(R)$  versus ellipticity  $\epsilon(R)$  for a typical S0 galaxy (top panel), and a typical E galaxy (bottom panel). Radius is indicated by the colour of each point. The ATLAS<sup>3D</sup> luminosity-weighted value for each galaxy is denoted by the large square. The corresponding  $v/\sigma(R)$  values, as calculated according to the scaling equation in Section 3, are indicated by the right axis. Arrows indicate the increasing radial direction for each track.

## 4 RESULTS

Before presenting the radial tracks of galaxies with intermediate-scale discs, we first display tracks typical of an E and an S0 galaxy. Fig. 1 shows the radial track of the lenticular galaxy NGC 1023, and the slow rotator elliptical galaxy NGC 4365. We note that while NGC 4365 has been observed to host a central kinematically distinct core (Krajnović et al. 2011), we do not resolve this feature, and therefore the larger radii kinematics depict those of a typical E galaxy.

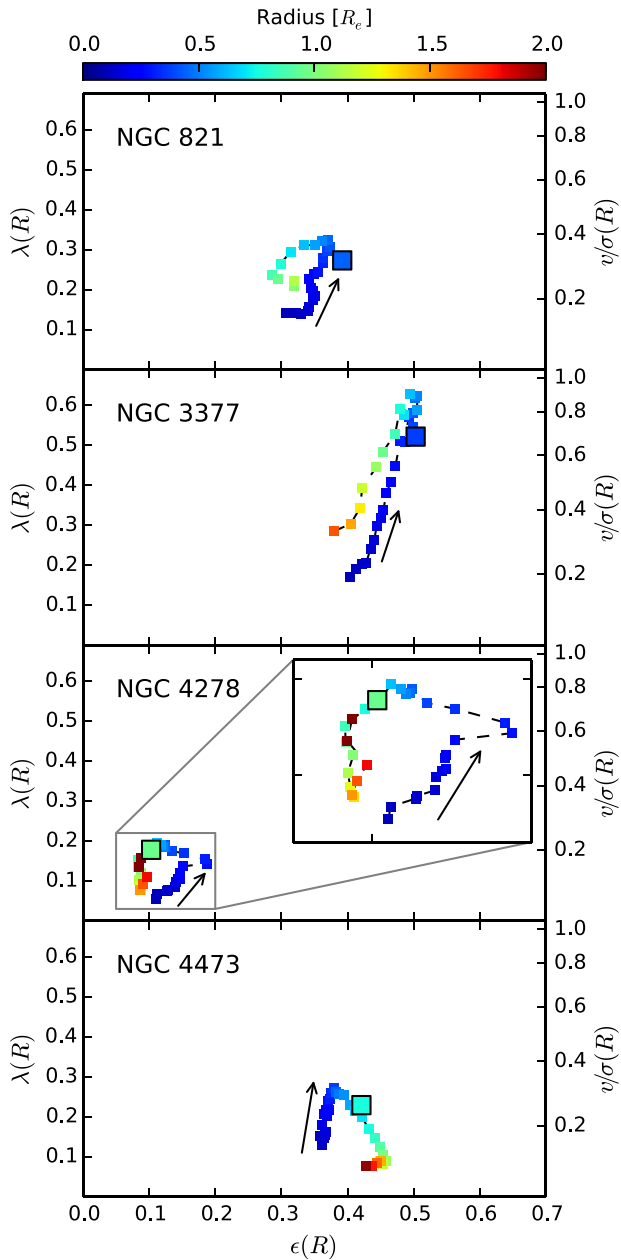
For NGC 1023, the ellipticity increases with radius, becoming more disc like in its outskirts, as expected for an S0 galaxy containing a large-scale disc that dominates the light at large radii. This increase in ellipticity is associated with a gradual increase in both  $\lambda(R)$  and  $v/\sigma(R)$ . At larger radii the rotation parameters and ellipticity increase more gradually than in the inner region, indicated by the radial bunching of the data towards the end of the track. Our data, and thus the track for NGC 1023, reaches a radial extent of  $\sim 2 R_e$ . The ‘global’  $\lambda_R - \epsilon$  value from ATLAS<sup>3D</sup> (Emsellem et al. 2011) is indicated on the plot as a large square. Like the points of the radial track, it is coloured according to the maximum radial extent of the data, i.e.  $0.40 R_e$  (which is why the ATLAS<sup>3D</sup> point is shown mid-way along the NGC 1023 radial track).

The track for NGC 4365 is much more stagnant. The ellipticity is roughly constant at  $0.2 < \epsilon < 0.3$  within  $\sim 1 R_e$ , and both  $\lambda(R)$  and  $v/\sigma(R)$  show little change with radius. The ATLAS<sup>3D</sup> data for this galaxy extend to  $0.19 R_e$ .

Moving on to our four declining spin galaxies, we show the radial tracks of NGC 821, NGC 3377, NGC 4278 and NGC 4473 in Fig. 2. For each of these galaxies, it can be seen that there is a downturn with increasing radius in both spin and ellipticity, which generally occurs at  $\sim 0.5 - 0.7 R_e$ . When the downturn of both properties occurs at the same radius, the track moves in an anticlockwise direction, as is the case for three out of the four galaxies. NGC 4473 is slightly different, in that the ellipticity downturn occurs at  $\sim 1 R_e$ , lagging that of the rotation which occurs at  $\sim 0.5 R_e$ . This results in a roughly clockwise track. This signature is due to the double-sigma nature of NGC 4473, which has a disc that is counter rotating with respect to a slowly rotating outer region (Krajnović et al. 2011). The counter rotating region has a diluting effect on the disc, which causes the  $\lambda(R)$  value to reduce inwards of the radius at which the local ellipticity transitions.

Fig. 3 displays tracks for 15 additional galaxies from the SLUGGS survey. The data for these tracks have previously been published in Arnold et al. (2014), Foster et al. (2016) and B17. Tracks for S0 galaxies move in a single direction towards the top right region of the plot, whereas tracks for E galaxies hover and remain in the bottom left region of the plot. The tracks for the three hybrid or ES galaxies from Fig. 2 have been plotted as blue lines, and although the endpoints of the tracks coincide with those of the E galaxies, the structure of the tracks themselves are quite different – particularly for NGC 3377, which strongly veers into the S0 region before turning down. The maximum radius for each track varies between  $1 - 3 R_e$ .

For clarity, we omit S0 galaxies with somewhat face-on discs (as identified in B17) from Fig. 3. Going from edge-on ( $i = 90^\circ$ ) to face-on ( $i = 0^\circ$ ), the measured rotation reduces from the intrinsic rotational velocity as  $v \sin i$ . The effect of this reduced velocity measurement is to lower the  $\lambda(R)$  and  $v/\sigma(R)$  values. Hence, S0 galaxies with both lower intrinsic rotation and lower inclination will reside in the lower left region of the modified diagram. The



**Figure 2.** Similar to Fig. 1 but for four galaxies whose local spin and ellipticity decline at larger radii. Large square points indicate the single aperture measurements from ATLAS<sup>3D</sup>, whose data extend to 0.44, 0.37, 0.95 and 0.76  $R_e$  for NGC 821, 3377, 4278 and 4473, respectively.

same effect is evident in the aperture-based  $\lambda_R - \epsilon$  diagram (see fig. 15 of B17).

## 5 DISCUSSION

While the tracks for the three ES galaxies in Fig. 2 are similar, each is unique in shape and size, revealing valuable insight into their host galaxy’s structure and dynamics. Furthermore, these tracks, plus that for the double-sigma galaxy, highlight the amount of information that is not conveyed by single spin and ellipticity values. These galaxies display radial tracks very different to those of either typical elliptical or lenticular galaxies. These differences

are not seen when summarizing the rotational behaviour of galaxies with a single aperture spin and ellipticity value. Indeed, all four galaxies are classified as (centrally) fast rotators (Emsellem et al. 2011).

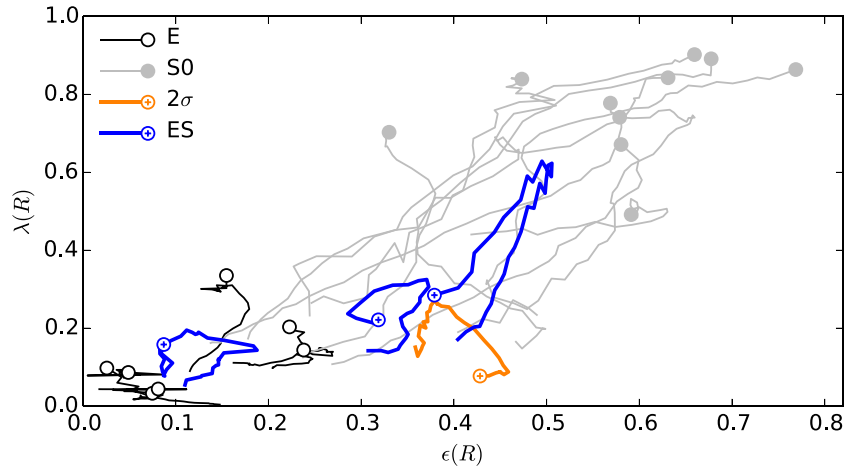
For each galaxy in Figs 1 and 2, we plot the  $\lambda_{R_e}$  value from ATLAS<sup>3D</sup> (Emsellem et al. 2011) as a large square. When using a larger aperture size, dramatic changes in the spin or ellipticity at greater radii will not strongly influence this value, since the flux-weighted measurement is driven by the rotationally supported galaxy centre (as portrayed in fig. 9 of B17). Hence, using larger apertures to measure  $\lambda_R$  values will not necessarily be sensitive to the changed behaviour. Since it is this behaviour at greater radii that differentiates the kinematics of galaxies with intermediate-scale discs from those of S0 galaxies, its capture is important.

Interestingly, NGC 4473 was one of the galaxies classified by Liller (1966) as an ‘ES’ galaxy, displaying characteristics mid-way between an elliptical and a lenticular galaxy, and was referred to as ‘ellicular’ by Graham et al. (2016). NGC 821, despite being previously classified as an E6 galaxy, was determined by Scorza & Bender (1995) to have a photometric stellar disc, while in the same study NGC 3377 was found to have an intermediate-scale disc within a boxy bulge. This disc within NGC 3377 was also identified by Arnold et al. (2014) from kinematics. Raimond et al. (1981) found an irregular, extended HI disc within NGC 4278 – a feature unusual in ‘elliptical’ galaxies. Hybrid elliptical/lenticular features have therefore been identified previously in each of these galaxies by varying methods, and the  $\lambda(R) - \epsilon(R)$  diagram portrays the hybrid nature clearly.

We expect that the slowly rotating, bulge-like outer regions of galaxies with intermediate-scale discs exhibit tracks similar to the outer regions of elliptical galaxies, with no strong features or variation. Hence, when probed far enough out, the ends of their tracks would likely converge to a fixed point in spin-ellipticity space, unless there is a change of ellipticity in the galaxy outskirts. Although the SLUGGS data have a relatively large radial extent of  $\sim 2 R_e$ , it is still not far enough to confirm this behaviour for the three galaxies with intermediate-scale discs. Discs can come in a range of sizes relative to the galaxy  $R_e$ , and as such their decline may or may not be detected by  $\sim 2 R_e$ . While NGC 4473 is not a typical spheroidal galaxy with an intermediate-scale disc (it is a double-sigma galaxy), the track does hover around a single point at larger radii, displaying elliptical-type behaviour in its outer regions.

The presence of intermediate-scale discs has been identified in some local ETGs, which have been interpreted as descendants of high- $z$  compact massive galaxies that have experienced small amounts of disc growth (Graham, Dullo & Savorgnan 2015; Graham et al. 2016; Savorgnan & Graham 2016). Spheroid masses for NGC 821, 3377 and 4473 were measured by Savorgnan et al. (2016), and determined to range from  $3.9 - 4.7 \times 10^{10} M_\odot$ . These spheroids are less massive than those of compact massive galaxies ( $\sim 10^{11} M_\odot$ ) studied by Graham et al. (2015), indicating that embedded, intermediate-scale discs occur within spheroids over a broad stellar mass range. Moreover, Graham et al. (2017) report a potential intermediate-scale disc in a dwarf ETG. If one assumes that the structural similarity between ES galaxies is indicative of a common formation path, then not only have the most massive ‘red nuggets’ at high- $z$  undergone some degree of disc formation to form present-day galaxies, but this process may apply to high- $z$  compact galaxies with a broad range of stellar masses.

In order to detect the presence and assess the prevalence of ES galaxies, it is necessary to analyse large galaxy samples. Several IFU



**Figure 3.** Tracks for SLUGGS galaxies in the modified spin-ellipticity diagram. Markers indicate the maximum radius end of each track.  $2\sigma$  refers to the double-sigma galaxy NGC 4473. It can be seen that the structure of the E and ES (NGC 821, 3377, 4278) galaxy tracks are very different. Low-inclination S0 galaxies (see B17 for details) have been omitted from this plot.

galaxy surveys, including Sydney AAO Multi-object Integral-field (Bryant et al. 2015), Mapping Nearby Galaxies at Apache Point Observatory (Bundy et al. 2015) and Calar Alto Legacy Integral Field Area (Sánchez et al. 2012) are in varying degrees of completion, and for a subsample of galaxies will reach sufficiently large radii to make such assessments. They will be able to make statistical statements about the prevalence and radial range of intermediate-scale discs in ETGs to build on the statistics provided by the ATLAS<sup>3D</sup> survey (Cappellari et al. 2011a), and give a clearer indication of the kinematic feature that separates, or more appropriately unites, elliptical and lenticular galaxies.

## 6 CONCLUSIONS

We show that the use of radial tracks in an annulus-based spin-ellipticity  $\lambda(R) - \epsilon(R)$  diagram is able to identify a population of discy elliptical early-type galaxies. This provides a significant advance over existing  $\lambda - \epsilon$  diagrams. Discy elliptical galaxies display fast rotation within their inner regions and slow rotation within their outer regions where their discs no longer dominate. In early applications of the spin-ellipticity diagram, which used a single, central flux-weighted value of spin and ellipticity, such galaxies could not be uniquely distinguished from other ETGs and required alternative means of identification. The  $\lambda(R) - \epsilon(R)$  diagram provides a succinct method of summarizing both kinematic and photometric data within a single diagram, facilitating a quick determination of the galaxy behaviour.

## ACKNOWLEDGEMENTS

The data presented herein were obtained at the *Spitzer* Space Telescope and the W. M. Keck Observatory. We acknowledge the work by Luciana Sinpetru in producing the ellipticity profiles used within this paper. SB acknowledges the support of the AAO PhD Topup Scholarship. AWG and DAF are supported under the ARC's funding schemes DP170102923 and DP130100388, respectively. JPB and AJR are supported by the NSF grants AST-16165598 and AST-1616710, respectively, AJR is supported as a Research Corporation for Science Advancement Cottrell Scholar and JS is supported by a Packard Fellowship.

## REFERENCES

- Alabi A. B. et al., 2015, MNRAS, 452, 2208  
 Andreon S., Davoust E., Michard R., Nieto J.-L., Poulain P., 1996, A&AS, 116, 429  
 Arnold J. A. et al., 2014, ApJ, 791, 80  
 Babcock H. W., 1939, Lick Obs. Bull., 19, 41  
 Bellstedt S., Forbes D. A., Foster C., Romanowsky A. J., Brodie J. P., Pastorello N., Alabi A., Villaume A., 2017, MNRAS, 467, 4540  
 Bertola F., Capaccioli M., 1975, ApJ, 200, 439  
 Binney J., 1978, MNRAS, 183, 501  
 Brodie J. P. et al., 2014, ApJ, 796, 52  
 Bryant J. J. et al., 2015, MNRAS, 447, 2857  
 Bundy K. et al., 2015, ApJ, 798, 7  
 Cappellari M. et al., 2011a, MNRAS, 413, 813  
 Cappellari M. et al., 2011b, MNRAS, 416, 1680  
 Cinzano P., van der Marel R. P., 1994, MNRAS, 270, 325  
 Davies R. L., 1981, MNRAS, 194, 879  
 Davies R. L., Efstathiou G., Fall S. M., Illingworth G., Schechter P. L., 1983, ApJ, 266, 41  
 de Zeeuw P. T. et al., 2002, MNRAS, 329, 513  
 di Tullio G., 1978, A&A, 62, L17  
 Emsellem E. et al., 2007, MNRAS, 379, 401  
 Emsellem E. et al., 2011, MNRAS, 414, 888  
 Forbes D. A., Sinpetru L., Savorgnan G., Romanowsky A. J., Usher C., Brodie J., 2017, MNRAS, 464, 4611  
 Foster C., Arnold J. A., Forbes D. A., Pastorello N., Romanowsky A. J., Spitler L. R., Strader J., Brodie J. P., 2013, MNRAS, 435, 3587  
 Foster C. et al., 2016, MNRAS, 457, 147  
 Graham A. W., Dullo B. T., Savorgnan G. A. D., 2015, ApJ, 804, 32  
 Graham A. W., Ciambur B. C., Savorgnan G. A. D., 2016, ApJ, 831, 132  
 Graham A. W., Janz J., Penny S. J., Chilingarian I. V., Ciambur B. C., Forbes D. A., Davies R. L., 2017, ApJ, 840, 68  
 Illingworth G., 1977, ApJ, 218, L43  
 Kormendy J., Illingworth G., 1982, ApJ, 256, 460  
 Krajnović D. et al., 2011, MNRAS, 414, 2923  
 Krajnović D. et al., 2013, MNRAS, 432, 1768  
 Larson R. B., 1975, MNRAS, 173, 671  
 Liller M. H., 1966, ApJ, 146, 28  
 Michard R., Marchal J., 1993, A&AS, 98, 29  
 Nieto J.-L., Capaccioli M., Held E. V., 1988, A&A, 195, L1  
 Pastorello N., Forbes D. A., Foster C., Brodie J. P., Usher C., Romanowsky A. J., Strader J., Arnold J. A., 2014, MNRAS, 442, 1003

Raimond E., Faber S. M., Gallagher J. S., III, Knapp G. R., 1981, *ApJ*, 246, 708

Raskutti S., Greene J. E., Murphy J. D., 2014, *ApJ*, 786, 23

Rix H.-W., White S. D. M., 1990, *ApJ*, 362, 52

Sánchez S. F. et al., 2012, *A&A*, 538, A8

Sandage A., Freeman K. C., Stokes N. R., 1970, *ApJ*, 160, 831

Savorgnan G. A. D., Graham A. W., 2016, *MNRAS*, 457, 320

Savorgnan G. A. D., Graham A. W., Marconi A., Sani E., 2016, *ApJ*, 817, 21

Schechter P. L., Gunn J. E., 1979, *ApJ*, 229, 472

Scorza C., Bender R., 1995, *A&A*, 293, 20

Simien F., Michard R., 1990, *A&A*, 227, 11

Williams T. B., 1975, *ApJ*, 199, 586

## APPENDIX: ELLIPTICITY PROFILES

The ellipticity profiles measured from *Spitzer* imaging (Forbes et al. 2017) are included here for each of the galaxies presented within this paper. The effective radii ( $R_e$ ) for each galaxy can be found in Forbes et al. (2017).

**Table A1.** NGC 821.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.10	0.308	0.24	0.343	0.56	0.351
0.11	0.322	0.26	0.305	0.63	0.334
0.12	0.329	0.29	0.355	0.70	0.314
0.14	0.340	0.32	0.363	0.78	0.299
0.15	0.342	0.35	0.363	0.87	0.286
0.16	0.348	0.38	0.369	0.95	0.295
0.18	0.351	0.42	0.373	1.02	0.319
0.20	0.347	0.46	0.371	1.13	0.319
0.22	0.345	0.51	0.362	–	–

**Table A2.** NGC 1023.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.10	0.178	0.30	0.285	0.80	0.483
0.11	0.180	0.33	0.303	0.86	0.508
0.12	0.183	0.35	0.339	0.92	0.535
0.14	0.192	0.38	0.345	0.98	0.566
0.15	0.199	0.42	0.354	1.05	0.589
0.16	0.211	0.45	0.376	1.13	0.605
0.18	0.221	0.49	0.399	1.21	0.624
0.19	0.228	0.53	0.420	1.31	0.635
0.21	0.235	0.57	0.439	1.44	0.638
0.23	0.246	0.62	0.453	1.57	0.643
0.25	0.258	0.68	0.465	1.72	0.649
0.28	0.271	0.71	0.510	1.86	0.659

**Table A3.** NGC 1407.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.10	0.044	0.22	0.044	0.42	0.036
0.11	0.044	0.24	0.047	0.46	0.044
0.12	0.041	0.26	0.048	0.51	0.061
0.13	0.039	0.29	0.046	0.56	0.042
0.15	0.040	0.32	0.053	0.63	0.005
0.16	0.042	0.35	0.053	0.66	0.087
0.18	0.048	0.38	0.048	0.74	0.049
0.20	0.044	–	–	–	–

**Table A4.** NGC 2549.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.10	0.244	0.30	0.442	0.98	0.510
0.11	0.244	0.33	0.463	1.07	0.525
0.12	0.234	0.36	0.469	1.15	0.544
0.14	0.227	0.39	0.470	1.23	0.565
0.15	0.233	0.44	0.467	1.33	0.581
0.16	0.252	0.48	0.459	1.45	0.590
0.18	0.290	0.53	0.453	1.58	0.595
0.19	0.314	0.59	0.447	1.74	0.594
0.21	0.337	0.65	0.445	1.91	0.600
0.22	0.362	0.71	0.450	2.08	0.607
0.24	0.380	0.78	0.457	2.26	0.615
0.26	0.396	0.84	0.474	2.44	0.631
0.28	0.418	0.91	0.497	–	–

**Table A5.** NGC 2768.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.10	0.269	0.26	0.406	0.67	0.531
0.11	0.281	0.29	0.421	0.73	0.542
0.12	0.297	0.31	0.436	0.79	0.553
0.13	0.309	0.34	0.450	0.86	0.563
0.14	0.323	0.37	0.468	0.93	0.576
0.16	0.332	0.40	0.480	1.01	0.588
0.17	0.342	0.43	0.493	1.10	0.599
0.19	0.353	0.47	0.502	1.20	0.603
0.21	0.364	0.51	0.510	1.33	0.597
0.22	0.377	0.56	0.513	1.46	0.599
0.24	0.390	0.61	0.522	1.64	0.580

**Table A6.** NGC 2974.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.11	0.237	0.24	0.329	0.54	0.372
0.12	0.245	0.26	0.344	0.60	0.354
0.13	0.259	0.28	0.352	0.67	0.338
0.14	0.281	0.31	0.359	0.74	0.334
0.15	0.304	0.34	0.359	0.82	0.328
0.16	0.309	0.37	0.363	0.90	0.326
0.18	0.315	0.41	0.370	0.99	0.327
0.20	0.324	0.45	0.376	1.09	0.326
0.22	0.324	0.49	0.380	1.20	0.330

**Table A7.** NGC 3115.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.10	0.405	0.31	0.574	0.95	0.652
0.11	0.413	0.34	0.574	1.04	0.655
0.12	0.431	0.37	0.573	1.14	0.657
0.13	0.450	0.41	0.569	1.25	0.660
0.15	0.464	0.45	0.571	1.39	0.653
0.16	0.480	0.49	0.576	1.54	0.651
0.17	0.491	0.53	0.584	1.71	0.642
0.19	0.505	0.58	0.592	1.90	0.637
0.20	0.520	0.63	0.606	2.10	0.633
0.22	0.532	0.68	0.615	2.34	0.621
0.24	0.544	0.74	0.623	2.61	0.611
0.26	0.556	0.81	0.631	2.94	0.592
0.28	0.565	0.88	0.640	3.29	0.579

**Table A8.** NGC 3608.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.11	0.162	0.25	0.179	0.56	0.232
0.12	0.165	0.27	0.179	0.61	0.240
0.13	0.171	0.30	0.181	0.68	0.235
0.14	0.173	0.33	0.183	0.75	0.232
0.15	0.177	0.36	0.185	0.82	0.239
0.17	0.178	0.39	0.190	0.91	0.226
0.18	0.180	0.43	0.203	1.00	0.220
0.20	0.180	0.47	0.214	1.10	0.223
0.22	0.179	0.51	0.226	–	–

**Table A9.** NGC 3377.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.10	0.404	0.25	0.466	0.63	0.495
0.11	0.413	0.27	0.472	0.70	0.485
0.12	0.422	0.30	0.480	0.77	0.479
0.13	0.428	0.33	0.486	0.86	0.472
0.15	0.435	0.36	0.493	0.96	0.454
0.16	0.440	0.39	0.498	1.07	0.444
0.17	0.444	0.43	0.500	1.19	0.422
0.19	0.450	0.47	0.502	1.32	0.419
0.21	0.454	0.52	0.506	1.47	0.405
0.23	0.459	0.57	0.505	1.65	0.379

**Table A10.** NGC 4111.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.10	0.489	0.42	0.499	1.39	0.618
0.11	0.480	0.46	0.495	1.52	0.624
0.12	0.479	0.52	0.479	1.66	0.629
0.14	0.465	0.57	0.473	1.81	0.634
0.16	0.443	0.63	0.476	1.97	0.641
0.17	0.436	0.69	0.482	2.13	0.653
0.19	0.435	0.75	0.493	2.30	0.668
0.21	0.442	0.81	0.507	2.47	0.684
0.23	0.447	0.88	0.523	2.64	0.701
0.24	0.465	0.95	0.541	2.82	0.717
0.27	0.466	1.02	0.561	3.05	0.727
0.29	0.469	1.10	0.578	3.27	0.740
0.33	0.467	1.18	0.594	3.53	0.750
0.36	0.469	1.28	0.610	3.74	0.769
0.39	0.487	–	–	–	–

**Table A11.** NGC 4278.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.10	0.110	0.28	0.188	0.85	0.083
0.11	0.112	0.31	0.184	0.93	0.083
0.12	0.127	0.34	0.152	1.02	0.088
0.13	0.128	0.38	0.134	1.12	0.085
0.15	0.139	0.42	0.125	1.24	0.086
0.16	0.140	0.47	0.123	1.36	0.088
0.18	0.143	0.51	0.121	1.50	0.087
0.20	0.146	0.57	0.118	1.64	0.090
0.22	0.146	0.63	0.112	1.80	0.097
0.24	0.146	0.69	0.105	1.99	0.084
0.26	0.152	0.76	0.095	2.19	0.087

**Table A12.** NGC 4365.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.10	0.205	0.23	0.248	0.54	0.263
0.11	0.204	0.25	0.247	0.59	0.264
0.12	0.209	0.28	0.246	0.65	0.263
0.13	0.218	0.31	0.231	0.71	0.27
0.14	0.227	0.34	0.239	0.78	0.268
0.16	0.236	0.37	0.254	0.86	0.269
0.17	0.241	0.41	0.259	0.96	0.248
0.19	0.243	0.45	0.254	1.07	0.238
0.21	0.244	0.49	0.256	–	–

**Table A13.** NGC 4374.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.10	0.147	0.23	0.114	0.45	0.067
0.11	0.144	0.25	0.109	0.50	0.063
0.13	0.141	0.27	0.104	0.55	0.051
0.14	0.141	0.30	0.093	0.60	0.060
0.15	0.136	0.33	0.089	0.67	0.037
0.17	0.136	0.37	0.080	0.74	0.042
0.18	0.126	0.41	0.079	0.82	0.026
0.20	0.120	–	–	–	–

**Table A14.** NGC 4473.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.11	0.358	0.31	0.374	0.82	0.441
0.12	0.361	0.34	0.375	0.90	0.449
0.13	0.365	0.37	0.379	0.99	0.453
0.14	0.366	0.40	0.382	1.08	0.459
0.16	0.368	0.44	0.382	1.19	0.454
0.17	0.362	0.49	0.387	1.32	0.450
0.19	0.364	0.53	0.395	1.45	0.450
0.21	0.370	0.58	0.403	1.60	0.445
0.23	0.367	0.64	0.411	1.77	0.438
0.25	0.372	0.69	0.422	1.96	0.428
0.28	0.371	0.75	0.433	2.17	0.422

**Table A15.** NGC 4486.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.10	0.018	0.23	0.054	0.50	0.044
0.11	0.021	0.26	0.047	0.55	0.044
0.12	0.031	0.28	0.049	0.60	0.059
0.13	0.062	0.31	0.037	0.66	0.057
0.14	0.094	0.34	0.033	0.72	0.065
0.15	0.111	0.38	0.032	0.79	0.069
0.17	0.090	0.42	0.035	0.87	0.072
0.19	0.075	0.46	0.039	0.96	0.081
0.21	0.077	–	–	–	–



**Table A16.** NGC 4494.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.11	0.115	0.27	0.153	0.71	0.150
0.12	0.116	0.30	0.162	0.78	0.155
0.13	0.117	0.33	0.170	0.86	0.144
0.14	0.119	0.36	0.175	0.94	0.147
0.16	0.126	0.39	0.179	1.04	0.143
0.17	0.133	0.43	0.178	1.14	0.149
0.19	0.138	0.48	0.173	1.26	0.139
0.21	0.145	0.53	0.169	1.40	0.127
0.23	0.149	0.58	0.161	1.51	0.155
0.25	0.150	0.64	0.152	1.66	0.155

**Table A17.** NGC 4526.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.10	0.467	0.38	0.373	1.16	0.504
0.11	0.473	0.42	0.361	1.26	0.524
0.12	0.489	0.46	0.360	1.35	0.545
0.13	0.496	0.51	0.362	1.44	0.569
0.14	0.499	0.56	0.360	1.54	0.596
0.16	0.489	0.61	0.372	1.64	0.619
0.18	0.474	0.66	0.387	1.78	0.632
0.20	0.461	0.72	0.407	1.92	0.647
0.22	0.450	0.78	0.427	2.07	0.658
0.25	0.435	0.84	0.447	2.24	0.669
0.27	0.418	0.91	0.461	2.44	0.677
0.31	0.401	0.99	0.474	2.68	0.677
0.34	0.386	1.08	0.488	2.95	0.677

**Table A18.** NGC 4564.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.11	0.138	0.34	0.281	0.86	0.536
0.12	0.140	0.37	0.301	0.94	0.549
0.13	0.148	0.40	0.322	1.02	0.560
0.14	0.153	0.44	0.345	1.11	0.568
0.16	0.157	0.47	0.368	1.21	0.573
0.17	0.171	0.51	0.390	1.33	0.577
0.19	0.183	0.55	0.414	1.46	0.578
0.20	0.196	0.59	0.438	1.60	0.581
0.22	0.205	0.64	0.463	1.76	0.581
0.24	0.215	0.68	0.486	1.94	0.578
0.27	0.230	0.74	0.504	2.16	0.569
0.29	0.243	0.80	0.523	2.38	0.569
0.32	0.261	–	–	–	–

**Table A19.** NGC 5846.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.10	0.062	0.20	0.059	0.39	0.062
0.11	0.062	0.22	0.057	0.43	0.067
0.13	0.062	0.25	0.055	0.47	0.079
0.14	0.064	0.27	0.054	0.52	0.078
0.15	0.068	0.30	0.059	0.57	0.082
0.17	0.071	0.33	0.058	0.63	0.086
0.18	0.062	0.36	0.061	0.69	0.075

**Table A20.** NGC 5866.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.21	0.419	0.42	0.592	1.00	0.593
0.23	0.439	0.45	0.613	1.10	0.590
0.25	0.457	0.49	0.623	1.21	0.591
0.27	0.474	0.54	0.626	1.32	0.598
0.29	0.490	0.59	0.630	1.45	0.599
0.32	0.512	0.65	0.630	1.60	0.598
0.34	0.531	0.72	0.623	1.76	0.598
0.37	0.552	0.81	0.609	1.94	0.593
0.39	0.575	0.90	0.600	2.14	0.591

**Table A21.** NGC 7457.

Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$	Radius ( $R_e$ )	$\epsilon$
0.10	0.246	0.27	0.399	0.65	0.471
0.11	0.260	0.29	0.410	0.68	0.509
0.12	0.281	0.32	0.413	0.77	0.481
0.13	0.306	0.35	0.417	0.86	0.467
0.15	0.308	0.38	0.422	0.94	0.473
0.16	0.316	0.42	0.420	1.05	0.457
0.17	0.336	0.46	0.421	1.15	0.464
0.19	0.345	0.50	0.444	1.25	0.477
0.21	0.352	0.54	0.447	1.39	0.462
0.22	0.371	0.60	0.444	1.52	0.473
0.24	0.387	–	–	–	–

This paper has been typeset from a  $\text{\TeX/L\AA\TeX}$  file prepared by the author.